## Black hole fireworks

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Until E & M all forces on a diagram come from contact







except for the force of gravity.

#### Action at a distance



That one body may act upon another at a distance through a vacuum without the mediation of anything else, by and through which their action and force may be conveyed from one another, is to me so great an absurdity that, I believe, no man who has in philosophic matters a competent faculty of thinking could ever fall into it.

Isaac Newton





## One What happens to collapsing matter?





# Two Toy model: the crystal ball

Three Exploding black holes



Small radii  $\rightsquigarrow$  deep quantum regime. Does an effective quantum pressure develop, avoiding a singularity?

Could this "pressure" push the matter back out? This would be like a cosmological bounce.

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho\left(1-\frac{\rho}{\rho_{\rm Pl}}\right)$$

![](_page_7_Picture_4.jpeg)

## Quantum effects

Quantum effects are relevant to black hole evolution.

In 1974 Stephen Hawking argued that black holes emit particles.

![](_page_8_Picture_3.jpeg)

His arguments can be cast as a quantum tunneling phenomenon.

Classically, only way to get over a barrier is with enough energy.

![](_page_9_Picture_1.jpeg)

Quantum mechanically you can go through barriers.

![](_page_10_Figure_1.jpeg)

However, the probabilities are small  $\rightsquigarrow$  long times.

![](_page_11_Figure_0.jpeg)

## Hawking radiation

![](_page_12_Figure_1.jpeg)

Vacuum is exciting—virtual pairs of particles are popping in and out of existence

They are virtual because one has +E and one -E

The -E particle is forbidden outside the horizon—tunneling inside it becomes allowed

The +E particle can escape to far away and carry some of the black hole's mass

## Hawking radiation

Because Hawking radiation is due to quantum tunneling we know that it must be slow. But how slow?

![](_page_13_Picture_2.jpeg)

## Hawking radiation

Because Hawking radiation is due to quantum tunneling we know that it must be slow. But how slow?

![](_page_14_Picture_2.jpeg)

Very, very slow  $T_H \sim M^3$ . For a solar mass black hole it takes  $T_H = 10^{75}$  secs. The age of the universe is  $T_U = 10^{17}$  secs.

## Bounce?

Hawking radiation focuses attention on the matter—what about the geometry?

![](_page_15_Picture_2.jpeg)

Our ideas:

- $E \text{ is conserved at } \infty \rightsquigarrow \\ elastic \text{ bounce}$
- Neglect Hawking radiation
- ♣ Quantum process ↔ tunneling of geometry Begins outside horizon
- GR is time reversal invariant—black to white hole bounce

How can these ideas be consistent with causality?

## One What happens to collapsing matter?

![](_page_16_Picture_1.jpeg)

![](_page_16_Picture_2.jpeg)

# Two Toy model: the crystal ball

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![](_page_16_Picture_5.jpeg)

## It's about time

![](_page_17_Figure_1.jpeg)

On the train:  $\Delta t_0 = \frac{h}{c}$ 

On the ground: 
$$\Delta t = \frac{\sqrt{h^2 + (v\Delta t)^2}}{c} \implies \Delta t = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \Delta t_0$$

"Moving clocks run slow!"

Because gravity bends *spacetime*, a clock lower in a gravitational potential runs slow!

How long does it take a shell of light to bounce off a mirrored ball?

$$\tau_R = \sqrt{1 - \frac{2M}{R}} \Big( R - a - 2M \ln \frac{a - 2M}{R - 2M} \Big)$$

![](_page_18_Picture_4.jpeg)

## How classical?

$$\tau_R = \sqrt{1 - \frac{2M}{R} \left( R - a - 2M \ln \frac{a - 2M}{R - 2M} \right)}$$

Classicality parameter

$$q = \ell_{\mathsf{PI}} \mathcal{R} \tau_R,$$

here  $\mathcal{R}\sim \frac{M}{R^3}$  is a measure of strength of curvature and q<<1 means classical.

q can be near 1 for  $a\sim 2M$  and  $\tau_R$  large enough. It has a maximum at  $R_q=\frac{7}{6}(2M)$  (outside horizon!) and requiring  $q\sim 1$  gives  $\tau_q\sim M^2$ .

![](_page_19_Picture_6.jpeg)

# Quantum gravity effects can take hold outside the horizon!

## One What happens to collapsing matter?

![](_page_21_Picture_1.jpeg)

![](_page_21_Picture_2.jpeg)

# Two Toy model: the crystal ball

Three Exploding black holes

![](_page_21_Picture_5.jpeg)

Let's try to build a solution of Einstein's equations where collapsing matter bounces back out.

• Idea: glue a black hole to a white hole.

![](_page_22_Figure_2.jpeg)

![](_page_22_Figure_3.jpeg)

In fact, a glued version of these two space times exists

![](_page_23_Figure_1.jpeg)

but it's upside down.

Can we cut it up?

#### Yes! Use the crossed fingers —

![](_page_24_Picture_1.jpeg)

![](_page_24_Picture_2.jpeg)

... and sew it all up...

#### The spacetime

![](_page_25_Figure_1.jpeg)

![](_page_26_Figure_0.jpeg)

#### Full metric: join the pieces

Spherical symmetry:

$$ds^{2} = -F(u,v)dudv + r^{2}(u,v)(d\theta^{2} + sin^{2}\theta d\phi^{2})$$

Region I (Flat): 
$$F(u_I,v_I)=1, \qquad r_I(u_I,v_I)=rac{v_I-u_I}{2}.$$
 Bounded by:  $v_I < 0.$ 

$$\begin{split} \text{Region II (Schwarzschild):} \quad F(u,v) &= \frac{32m^3}{r}e^{\frac{r}{2m}} \qquad \left(1-\frac{r}{2m}\right)e^{\frac{r}{2m}} = uv. \\ \text{Matching:} \quad r_I(u_I,v_I) &= r(u,v) \quad \longrightarrow \quad u(u_I) = \frac{1}{v_o}\left(1+\frac{u_I}{4m}\right)e^{\frac{u_I}{4m}}. \\ \text{Region III (Quantum):} \qquad F(u_q,v_q) &= \frac{32m^3}{r_q}e^{\frac{r_q}{2m}}, \qquad r_q = v_q - u_q. \end{split}$$

- Collapsing matter bounces in a short time locally but a long time from far away,  $\sim M^2$ . Solar mass:  $\tau_q \sim 10^{32}$  sec,  $\tau_H \sim 10^{75}$  sec,  $\tau_U \sim 10^{17}$  sec.
- Possible to describe using a metric with no singularity, two trapped regions, and all matter exiting ~> all info escapes
- Could a black hole be a bouncing star seen in super slow motion? With the constructed metric we can attack this question rigorously.
- I want to calculate the WKB amplitude for a gravitational instanton giving this bounce process; now I can in principle!