

$$\dot{\vec{P}} = M \dot{\vec{R}} \quad \left(\begin{array}{l} \vec{R} = \text{CM position} \\ M = \text{total mass} \end{array} \right)$$

$$\dot{\vec{P}} = \vec{F}_{\text{ext}} = M \ddot{\vec{R}}$$

$$\dot{\vec{L}} = \dot{\vec{R}} \times \dot{\vec{P}} + \sum_{\alpha} \vec{r}'_{\alpha} \times M_{\alpha} \dot{\vec{v}}'_{\alpha}$$

position relative to CM ($\vec{r}'_{\alpha} = \vec{R} + \vec{r}'_{\alpha}$)

$$\dot{\vec{L}} = \dot{\vec{L}}(\text{motion of CM}) + \dot{\vec{L}}(\text{motion relative to CM}).$$

$$\frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} M_{\alpha} \dot{\vec{r}}_{\alpha}^2$$

Einstein summation

or

$$T = T(\text{motion of CM}) + T(\text{motion relative to CM})$$

of for a rigid body

$$T = T(\text{motion of CM}) + T(\text{rotation about relative to CM})$$

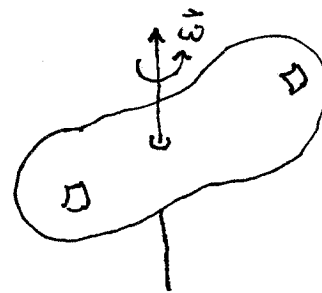
and

$$U = U_{\text{ext}} + U_{\text{int}}$$

for rigid body w/ conservative internal forces, a constant that can be dropped.

III Why must we distinguish $\vec{\omega}$ and \vec{L} ?

Example: Fixed axis of rotation.



"spared potato"

Choose fixed axis to be z-axis:

$$\vec{\omega} = (0, 0, \omega)$$

We want to calculate

$$\vec{L} = \sum \vec{l}_{\alpha} = \sum \vec{r}_{\alpha} \times M_{\alpha} \vec{v}_{\alpha}$$

Well,

$$\vec{r}_\alpha = (x_\alpha, y_\alpha, z_\alpha)$$

$$\text{and } \vec{v}_\alpha = (\vec{\omega} \times \vec{r}_\alpha) \quad (\text{recall } \frac{d\vec{r}}{dt} = \vec{\omega} \times \vec{r})$$

$$\vec{v}_\alpha = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0 & \omega \\ x_\alpha & y_\alpha & z_\alpha \end{vmatrix} = (-y_\alpha \omega, x_\alpha \omega, 0)$$

$$L_\alpha = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ x_\alpha & y_\alpha & z_\alpha \\ -m_\alpha y_\alpha \omega & m_\alpha x_\alpha \omega & 0 \end{vmatrix} = m_\alpha \omega (-x_\alpha z_\alpha, -y_\alpha z_\alpha, x_\alpha^2 + y_\alpha^2) \quad (\text{no sum})$$

So, why must we distinguish $\vec{\omega}$ and \vec{L} ? [The answer is even richer than why we distinguish \vec{v} and \vec{p} !] Unlike \vec{v} and \vec{p} they don't just differ by a scalar ($\vec{p} = m\vec{v}$) but they can even point in different directions (!!!):

$\vec{\omega}$ is not always parallel to \vec{L} .

Then

$$L_z = \sum m_\alpha (x_\alpha^2 + y_\alpha^2) \omega = m_\alpha r_\alpha^2 \omega = I_z \omega$$

with $I_z = m_\alpha r_\alpha^2$. This may look familiar. But

$$L_x = - \sum m_\alpha x_\alpha z_\alpha \omega$$

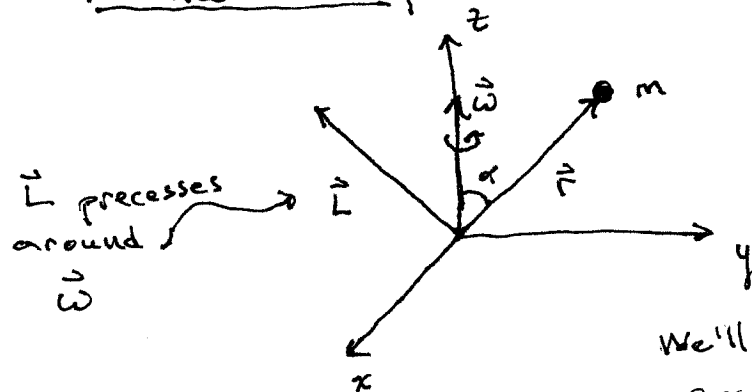
$$L_y = - \sum m_\alpha y_\alpha z_\alpha \omega$$

are new.

These prefactors are called the "products of inertia"

Why? \vec{L} depends on how the mass is distributed around the rotation axis, that is, on the body's shape.

Another example:



We'll return to this example.

We've done two examples, now let's do the general case.

IV Inertia Tensor

General case has

$$\vec{\omega} = (\omega_x, \omega_y, \omega_z)$$

and

$$\begin{aligned} \vec{L} &= \sum_{\alpha} m_{\alpha} \vec{r}_{\alpha} \times \vec{v}_{\alpha} \\ &= \sum_{\alpha} m_{\alpha} \vec{r}_{\alpha} \times (\vec{\omega} \times \vec{r}_{\alpha}) \end{aligned}$$

Recall (or use supplement on indices $\frac{2}{3}$ to show) the

Thus, after putting in $\sum_{\alpha} m_{\alpha}$,

we have

$$\begin{pmatrix} L_x \\ L_y \\ L_z \end{pmatrix} = \begin{pmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{yx} & I_{yy} & I_{yz} \\ I_{zx} & I_{zy} & I_{zz} \end{pmatrix} \begin{pmatrix} \omega_x \\ \omega_y \\ \omega_z \end{pmatrix}$$

non-diag contains "products of inertia"

diag contains "moments of inertia"

Reading off of (E_B*) we have

$$I_{xx} = \sum_{\alpha} m_{\alpha} (y_{\alpha}^2 + z_{\alpha}^2) \text{ etc.}$$

and

$$I_{xy} = - \sum_{\alpha} m_{\alpha} x_{\alpha} y_{\alpha} \text{ etc.}$$

bac - cab rule:

$$\vec{a} \times (\vec{b} \times \vec{c}) = \vec{b}(\vec{a} \cdot \vec{c}) - \vec{c}(\vec{a} \cdot \vec{b})$$

So that,

$$\begin{aligned} \vec{r} \times (\vec{\omega} \times \vec{r}) &= \vec{\omega}(r^2) - \vec{r}(\vec{r} \cdot \vec{\omega}) \\ &= (\omega_x r^2 - x(x\omega_x + y\omega_y + z\omega_z), \\ &\quad \omega_y r^2 - y(x\omega_x + y\omega_y + z\omega_z), \\ &\quad \omega_z r^2 - z(x\omega_x + y\omega_y + z\omega_z)) \\ &= (\omega_x(y^2 + z^2) - xy\omega_y - xz\omega_z, \\ &\quad \omega_y(x^2 + z^2) - yx\omega_x - yz\omega_z, \quad (E_B^*) \\ &\quad \omega_z(x^2 + y^2) - zx\omega_x - zy\omega_y) \end{aligned}$$

Shorthand notation

$$L_i = I_{ij} \omega_j$$

Einstein summation on j index

or

$$\vec{L} = \underline{\underline{I}} \vec{\omega}$$

$\underline{\underline{I}}$ is the moment of inertia or inertia tensor. For practical purposes

its just a matrix. While we have it up, note that $I_{xy} = I_{yx}$ etc, or

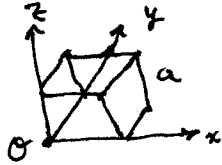
$$I_{ij} = I_{ji}$$

or in matrix notation

$$\underline{\underline{I}} = \underline{\underline{I}}^T \leftarrow \text{transpose} = \text{flip matrix along its main diagonal.}$$

Example: Cube: continuous

$$\rho = M/a^3$$



$$\begin{aligned} I_{xx} &= \int_0^a \int_0^a \int_0^a dx dy dz \rho (y^2 + z^2) \\ &= \rho \left(a^2 \left(\frac{y^3}{3} \Big|_0^a + a^2 \left(\frac{z^3}{3} \Big|_0^a \right) \right) \right) \\ &= \frac{2}{3} \rho a^5 = \frac{2}{3} M a^2 \end{aligned}$$

Next time: We'll see how

$$\underline{\underline{I}} = \underline{\underline{I}}^T$$

can significantly simplify things.

$$\begin{aligned} I_{xy} &= - \int_0^a \int_0^a \int_0^a \rho xy dx dy dz \\ &= -\rho a \left(\frac{x^2}{2} \Big|_0^a \left(\frac{y^2}{2} \Big|_0^a \right) \right) \\ &= -\frac{1}{4} \rho a^5 = -\frac{1}{4} M a^2 \end{aligned}$$

By symmetry we have

$$I = M a^2 \begin{pmatrix} \frac{2}{3} & -\frac{1}{4} & -\frac{1}{4} \\ -\frac{1}{4} & \frac{2}{3} & -\frac{1}{4} \\ -\frac{1}{4} & -\frac{1}{4} & \frac{2}{3} \end{pmatrix} \quad [\text{about corner}]$$