

Today's Outline: Lecture 3
Aug. 31st, 2011

I last lecture

Homogeneous linear Diff. Eqs

II last lecture &

damped oscillator wrap up

III Driven damped oscillator

$a_n \frac{d^n x}{dt^n} + \dots + a_1 \frac{dx}{dt} + a_0 x = 0$. Any equation of this form can be solved by a solution of the form:

$$x(t) = e^{rt}$$

with r constant (our standard guess).

Why it works: $x = e^{rt} \Rightarrow \frac{dx}{dt} = r e^{rt}$
 $\frac{d^2x}{dt^2} = r^2 e^{rt}, \dots, \frac{d^n x}{dt^n} = r^n e^{rt}$

Homogeneous linear eqns w/ P/C
constant coeff.:

Linear: $a_n(t) \frac{d^n x}{dt^n} + a_{n-1}(t) \frac{d^{n-1} x}{dt^{n-1}} + \dots + a_1(t) \frac{dx}{dt} + a_0(t) x = f(t)$

Homogeneous: $f(t) = 0$; constant coeff.:
 ~~$a_0(t), a_1(t)$~~ $a_0, a_1, \dots, a_{n-1}, a_n$ do not depend on t.

$$a_n r^n e^{rt} + a_{n-1} r^{n-1} e^{rt} + \dots + a_0 e^{rt} = 0$$

$$\Rightarrow a_n r^n + a_{n-1} r^{n-1} + \dots + a_1 r + a_0 = 0$$

Fundamental Theorem of Algebra guarantees that there are n solutions to this equation.

Clean up: Collect the a_i and the derivatives into a differential operator:

$$D = a_n \frac{d^n}{dt^n} + a_{n-1} \frac{d^{n-1}}{dt^{n-1}} + \dots + a_1 \frac{d}{dt} + a_0$$

Example: damped oscillator

$$D = \frac{d^2}{dt^2} + 2\beta \frac{d}{dt} + \omega_0^2$$

$$Dx = \ddot{x} + 2\beta \dot{x} + \omega_0^2 x$$

This operator is linear

$$D(ax) = a Dx \quad D(x_1 + x_2) = Dx_1 + Dx_2$$

For n^{th} order equations the linear combination with n constants is the general solution

$$x(t) = C_1 e^{r_1 t} + C_2 e^{r_2 t} + \dots + C_n e^{r_n t}$$

with r_i solutions of (*).

Except: if the algebraic equation gives multiple roots, then I'm missing solutions.

P2/6

Last time we recognized that the damped oscillator EoM could be solved by the standard guess, and found

For linear/homogeneous diff. eq.s any linear combination of solutions is a solution because

$$\begin{aligned} D(ax_1 + bx_2) &= a Dx_1 + b Dx_2 \\ &= 0 + 0 = 0. \checkmark \end{aligned}$$

In this case the "extra" solution is $t e^{r_i t}$. In general, for multiple roots (repeated m times)

$$e^{r_i t}, t e^{r_i t}, t^2 e^{r_i t}, \dots, t^{m-1} e^{r_i t}$$

Last time we recognized that the damped oscillator EoM could be solved by the standard guess, and found

$$x(t) = C_1 e^{r_1 t} + C_2 e^{r_2 t}$$

with $r_1 = -\beta + \sqrt{\beta^2 - \omega_0^2}$ $r_2 = -\beta - \sqrt{\beta^2 - \omega_0^2}$
damping coeff.

$$\beta = \frac{b}{2m} \quad \omega_0 = \sqrt{\frac{k}{m}} \quad B = \sqrt{\beta^2 - \omega_0^2}$$

$$\omega_1 = \sqrt{\omega_0^2 - \beta^2}$$

Started to explore three regimes

(1) Overdamped: $\beta > \omega_0$. Exponentially damped
 with two characteristic times $\tau_1 = \frac{1}{-\beta + B}$
 $\tau_2 = \frac{1}{\beta + B}$:

Also two times at work here:

damping $\gamma_1 = \gamma_2 = \beta$; and the period
 of the oscillations

$$\tau_2 = \frac{2\pi}{\omega_1}$$

Question: How many oscillations
 in one characteristic time?

$$\# = \frac{\tau_1}{\tau_2} = \frac{1}{\beta \cdot \frac{2\pi}{\omega_1}} = \frac{1}{\pi} \frac{\omega_1}{2\beta} \approx \frac{1}{\pi} \frac{\omega_0}{2\beta}$$

large mass
small damping or

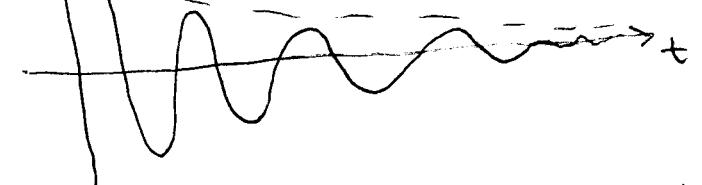
$$Q(\text{quality factor}) = \frac{\omega_0}{2\beta}$$

(2) Under-damped: $\omega_0 > \beta$ P3/6

Showed that

$$x(t) = A e^{-\beta t} \cos(\omega_1 t - \delta)$$

$x(t)$
exp. envelope
 $A e^{-\beta t}$



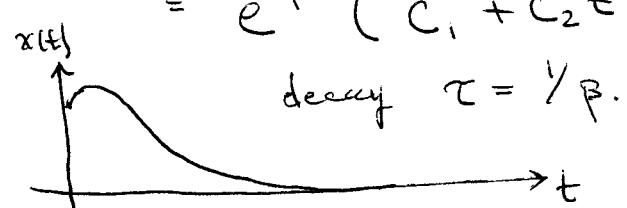
Characteristic time for
 damping is $\tau_1 = 1/\beta$.

Tells you how many times
 it rings.

(3) Critical damping: $\beta = \omega_0$
 $\Rightarrow r_1 = r_2$ multiple roots!

General solution is,

$$\begin{aligned} x(t) &= C_1 e^{-\beta t} + C_2 t e^{-\beta t} \\ &= e^{-\beta t} (C_1 + C_2 t) \end{aligned}$$



That's it for damping now we add in driving.

III Driven damped oscillator:

Think of a child on a swing oscillator (swing), friction (damping), driver (you pushing the child).

$$F_{\text{net}} = -kx - bx + F(t)$$

In general driving depends on time.

ω_0 the "natural frequency" and ω the "driving frequency".

This is an example of an ^{— stopped} here inhomogeneous ^{linear} diff. eq.

Suppose $x_h(t)$ is the general solution to the associated homogeneous equation then $\mathcal{D}^{\text{diff operator}} x_h = 0$

$$m\ddot{x} + bx + bu = F(t) \quad P4/6$$

$$\Rightarrow \ddot{x} + 2\beta\dot{x} + \omega_0^2 x = f(t) \equiv \frac{F(t)}{m}$$

f is the force per unit mass,

We'll be particularly interested in $f(t) = f_0 \cos(\omega t)$

('cause it's physically reasonable and solvable — can build more complex driving out of ~~one~~ cosines too)
Note well the distinction btwn

$$\mathcal{D} = \frac{d^2}{dt^2} + 2\beta \frac{d}{dt} + \omega_0^2$$

Suppose further that $x_p(t)$ is any old ("particular") solution to the inhomog. eq.

$$\mathcal{D} x_p(t) = f(t).$$

Then $x_h + x_p$ solves the inhomog. eq

$$\mathcal{D}(x_h + x_p) = \mathcal{D}x_h + \mathcal{D}x_p = 0 + f(t)$$

and is, in fact, the general solution.

[Note: It doesn't matter what particular solution you use (!)]

- the difference can be sealed up in the constants in x_h .

That is, $x_{p_1} - x_{p_2}$, is a solution of the homog. eq:

$$\begin{aligned} &= D(x_{p_1}) - D(x_{p_2}) \\ D(x_{p_1} - x_{p_2}) &= f(t) - f(t) \\ &= 0 \end{aligned}$$

The standard guess has shown how nice complex exponentials are. So we view this EoM as the real part of

$$\ddot{z} + 2\beta\dot{z} + \omega_0^2 z = f_0 e^{i\omega t}$$

with $z(t) = x(t) + iy(t)$. The imaginary part is then

$$\ddot{y} + 2\beta\dot{y} + \omega_0^2 y = f_0 \sin \omega t.$$

So, to solve driven P5/b case we need to focus on finding a particular solution to the inhomoq.

e.g. Do this by hook or by crook!

Want to solve

$$\ddot{x} + 2\beta\dot{x} + \omega_0^2 x = f_0 \cos(\omega t)$$

Try to guess a particular solution

$$z_p = C e^{i\omega t}$$

Physically this says that if we drive an oscillator at frequency ω we ^{might} expect it to respond at the same frequency.

Putting this,

$$\ddot{z}_p = -C\omega^2 e^{i\omega t} \quad \text{and} \quad \dot{z}_p = i\omega C e^{i\omega t}$$

and

$$-C\omega^2 e^{j\omega t} + 2\beta i \omega C e^{j\omega t} + \omega_0^2 C e^{j\omega t} = f_0 e^{j\omega t}$$

$$\Rightarrow C = \frac{f_0}{\omega_0^2 - \omega^2 + 2\beta i \omega}$$

That gives us a particular solution!

We will work on making Z_p

look simpler next time.