

Today's Outline

- I Last Step of Euler's Method
- II Example of Euler's method
- III Lagrange's method in the calculus of variations
- IV Examples: Geodesics on the cylinder

Lecture 6

Sept. 9th, 2011

I.

Last time we arrived at $P1/6$ the equation

$$\frac{\partial f}{\partial y_n} - \frac{1}{\Delta x} \left(\frac{\partial f}{\partial y'_n} - \frac{\partial f}{\partial y'_n} \right) = 0$$

as a condition for the sum

$$\sum_0^a f(x, y, y') \Delta x$$

with respect to variations of y_n .

to be extremized, See Figure 1 for notation.

In the limit as $\Delta x \rightarrow 0$ this condition becomes

$$\frac{\partial f}{\partial y(x)} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'(x)} \right) = 0$$

Notice that our derivation holds for any point of the curve, that is an arbitrary x and $y(x)$, hence all of them. Thus in using Euler's method you can always focus on a small segment of the curve.

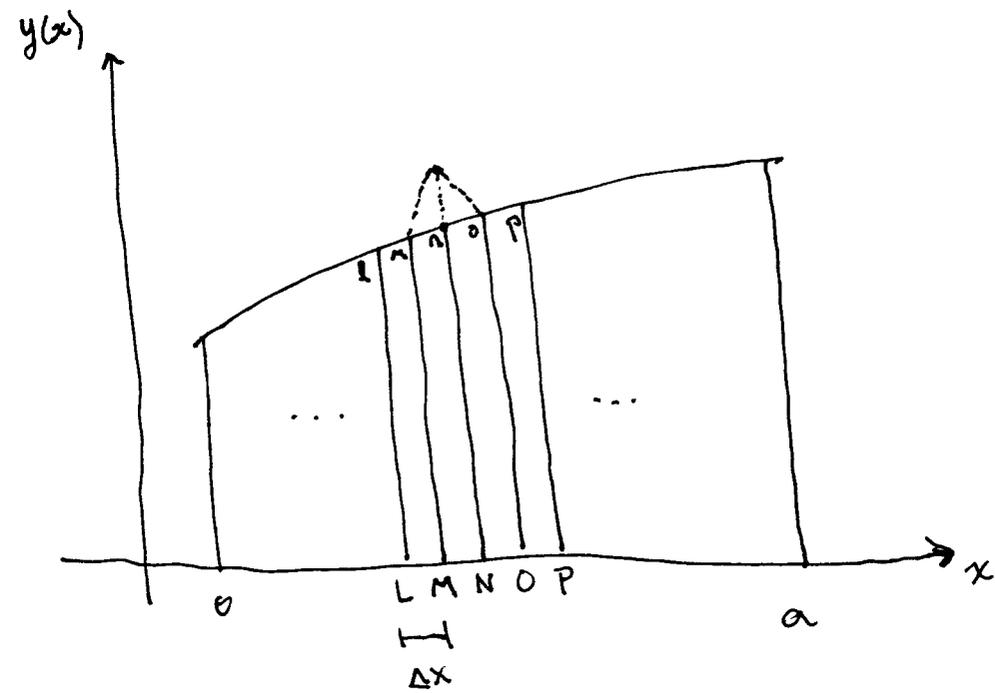
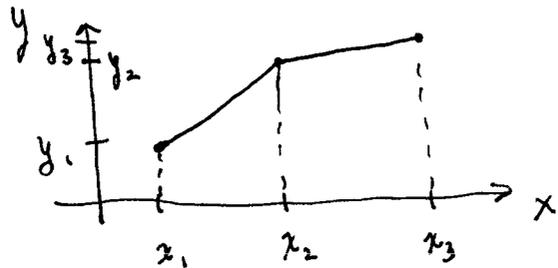


Figure 1.

II Let's do an example to illustrate this.

What is the shortest path connecting two points in the xy -plane?

Focus first on two nearby point (x_1, y_1) and (x_3, y_3)



So we impose ~~calculate~~,
calculate,

$$\frac{\partial L}{\partial y_2} = 0$$

$$\Rightarrow \frac{1}{2} \frac{1}{\sqrt{\Delta x^2 + (y_2 - y_1)^2}} \cdot 2(y_2 - y_1) + \frac{1}{2} \frac{1}{\sqrt{\Delta x^2 + (y_3 - y_2)^2}}$$

$$\times 2(y_3 - y_2) \cdot (-1) = 0$$

$$\Rightarrow \frac{y_2 - y_1}{\sqrt{\Delta x^2 + (y_2 - y_1)^2}} = \frac{y_3 - y_2}{\sqrt{\Delta x^2 + (y_3 - y_2)^2}}$$

The length of the path connecting them depends on how we choose the intermediate point (x_2, y_2) .

This length is just

$$L = \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2} + \sqrt{(x_3 - x_2)^2 + (y_3 - y_2)^2}$$
$$= \sqrt{\Delta x^2 + (y_2 - y_1)^2} + \sqrt{\Delta x^2 + (y_3 - y_2)^2}$$

We want to choose y_2 such that L is a minimum.

Factor out Δx^2 , and

$$\frac{(y_2 - y_1)/\Delta x}{\sqrt{1 - \frac{(y_2 - y_1)^2}{\Delta x^2}}} = \frac{(y_3 - y_2)/\Delta x}{\sqrt{1 - \frac{(y_3 - y_2)^2}{\Delta x^2}}}$$

and we notice these expressions are equal if

$$\left(\frac{y_2 - y_1}{\Delta x} \right) = \left(\frac{y_3 - y_2}{\Delta x} \right),$$

that is,

the slope is constant!
It's a straight line.

Imposing this condition, constant slope, on each segment of the curve gives a straight line for the whole curve.

III Lagrange's Method in the

calculus of variations
 Want to find the curve $y(x)$ that extremizes

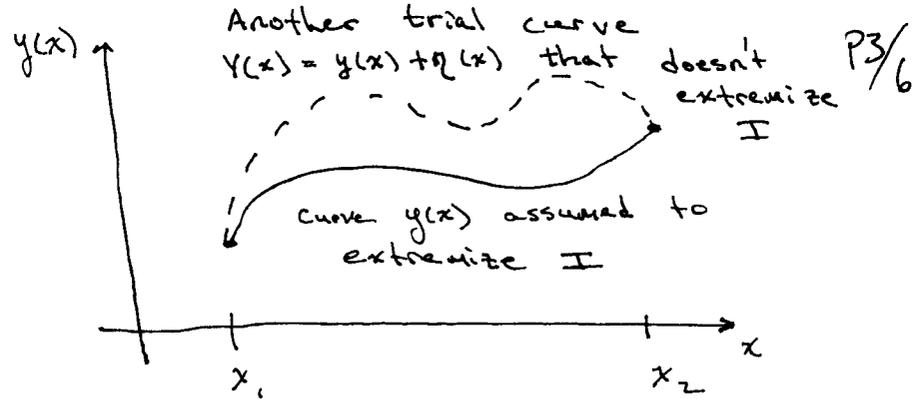
$$I = \int_{x_1}^{x_2} f(x, y(x), y'(x)) dx$$

Write your trial solution

$$Y(x) = y(x) + \alpha \eta(x)$$

in terms of a parameter α . Now you can turn on your variation of the whole curve just by changing α . In particular I becomes a normal function of α , $I(\alpha)$. So the condition that I is extremized is just

$$\frac{dI}{d\alpha} = 0$$



Lagrange also wants to turn this into ~~another~~ a calculus problem. He found a very clever way to do this:

[Subtlety: We assume $y(x)$ is the "right" curve so that we know where to evaluate the derivative, namely,

$$\left. \frac{dI}{d\alpha} \right|_{\alpha=0} = 0.]$$

Lets do the calculation.

If $Y(x) = y(x) + \alpha \eta(x)$ then

$$Y'(x) = y'(x) + \alpha \eta'(x).$$

So,

$$I(\alpha) = \int_{x_1}^{x_2} f(x, y, y') dx$$

$$= \int_{x_1}^{x_2} f(x, y(x) + \alpha \eta, y' + \alpha \eta') dx$$

Then,

$$\begin{aligned} \frac{dI}{d\alpha} &= \frac{d}{d\alpha} \int_{x_1}^{x_2} f dx \\ &= \int_{x_1}^{x_2} \frac{\partial f}{\partial \alpha} dx \end{aligned}$$

taking the derivative inside the integral makes it a partial.

To simplify the second term we use integration by parts

$$\int_{x_1}^{x_2} \frac{d}{dx} (f(x)g(x)) dx = f(x)g(x) \Big|_{x_1}^{x_2}$$

↳ The integral is the antiderivative.

But also

$$\int_{x_1}^{x_2} \frac{d}{dx} (fg) dx = \int_{x_1}^{x_2} f'g dx + \int_{x_1}^{x_2} fg' dx$$

Putting these together

$$\int_{x_1}^{x_2} \frac{df}{dx} g dx = f(x)g(x) \Big|_{x_1}^{x_2} - \int_{x_1}^{x_2} f \frac{dg}{dx} dx$$

$$\frac{\partial}{\partial \alpha} [f(x, y + \alpha \eta, y' + \alpha \eta')] \quad \text{P4/6}$$

$$\begin{aligned} &= \frac{\partial f}{\partial y} \cdot \frac{\partial}{\partial \alpha} (y + \alpha \eta) + \frac{\partial f}{\partial y'} \cdot \frac{\partial}{\partial \alpha} (y' + \alpha \eta') \\ \text{chain rule} &= \frac{\partial f}{\partial y} \eta + \frac{\partial f}{\partial y'} \eta' \end{aligned}$$

$$\text{So, } \frac{dI}{d\alpha} = \int_{x_1}^{x_2} \left(\frac{\partial f}{\partial y} \eta + \frac{\partial f}{\partial y'} \eta' \right) dx$$

"You can switch the derivative at the cost of a minus sign and a boundary term."

Then

$$\int_{x_1}^{x_2} \frac{\partial f}{\partial y'} \frac{d\eta}{dx} dx = \underbrace{\left(\frac{\partial f}{\partial y'} \eta \right) \Big|_{x_1}^{x_2}}_{=0} - \int_{x_1}^{x_2} \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \eta dx$$

and

↳ because of assumed bndry conditions.

$$\frac{dI}{d\alpha} = \int_{x_1}^{x_2} \left(\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \right) \eta dx$$

Want to set

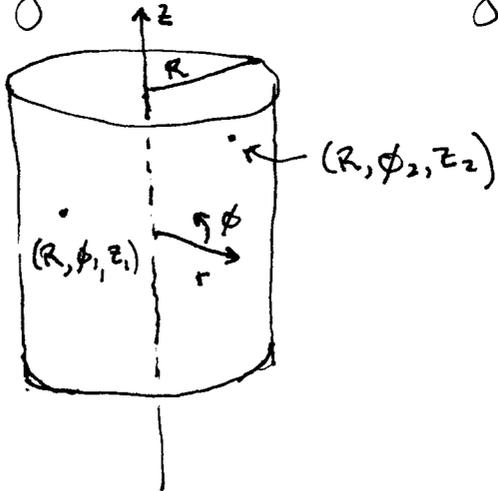
$$\frac{dI}{dy} = 0.$$

The pinnacle of the α cleverness is that this equation has to hold for all η . So what can we say when

$$\int_{x_1}^{x_2} g(x) \eta(x) dx = 0$$

for all η ? Well suppose $g \neq 0$ and choose η to be positive whenever g is and similarly for negative

Find the geodesics on a cylinder.



Describe as $\phi(z)$.

values. Then $g(x)\eta(x) > 0$ η always P5/6
and $\int_{x_1}^{x_2} g(x)\eta(x) dx \neq 0$

a contradiction! $\Rightarrow \boxed{g(x) = 0}$

So,
$$\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0 \quad \text{E-L Eq. 3 again!}$$

That completes our second proof.

Can go to examples (at last!).

$$ds^2 = R^2 d\phi^2 + dz^2$$

$$\Rightarrow ds = \sqrt{R^2 \left(\frac{d\phi}{dz} \right)^2 + 1} dz$$

$$L = \int_{z_1}^{z_2} \underbrace{\sqrt{R^2 (\phi')^2 + 1}} dz$$

$$f(z, \phi, \phi') = f(\phi')$$

E-L eq.

$$\frac{\partial f}{\partial \phi} - \frac{d}{dz} \left(\frac{\partial f}{\partial \phi'} \right) = - \frac{d}{dz} \left(\frac{\partial f}{\partial \phi'} \right) = 0$$

$$\Rightarrow \frac{\partial f}{\partial \phi'} = \text{const.} \stackrel{\text{call it}}{=} C_1$$

$$\Rightarrow \frac{1}{\sqrt{R^2 \phi'^2 + 1}} \cdot 2R^2 \phi' = C_1$$

$$\Rightarrow R^4 \phi'^2 = C_1^2 (R^2 \phi'^2 + 1)$$

$$\Rightarrow (R^4 - C_1^2 R^2) \phi'^2 = C_1^2$$

$$\Rightarrow \phi' = \text{const.} \stackrel{\text{call it}}{=} C$$

Then

$$\phi(z) = Cz + \text{constant } K$$

$$\left. \begin{aligned} \phi(z_1) &= Cz_1 + K = \phi_1 \\ \phi(z_2) &= Cz_2 + K = \phi_2 \end{aligned} \right\} \text{ solve } C \text{ and } K$$

This is the equation of a helix:

