

Today

I Lab Analysis

II Recollection of

Day 14 (before Spring break)

Time Examined

Day 16

I Reviewed how to do the Lab analysis.

II We studied how to

describe an oscillator mathematically

P1/3

Ideal oscillator:  $\Phi(t)$  the total phase

$$V(t) = V_0 \cos(2\pi f_0 t + S_0)$$

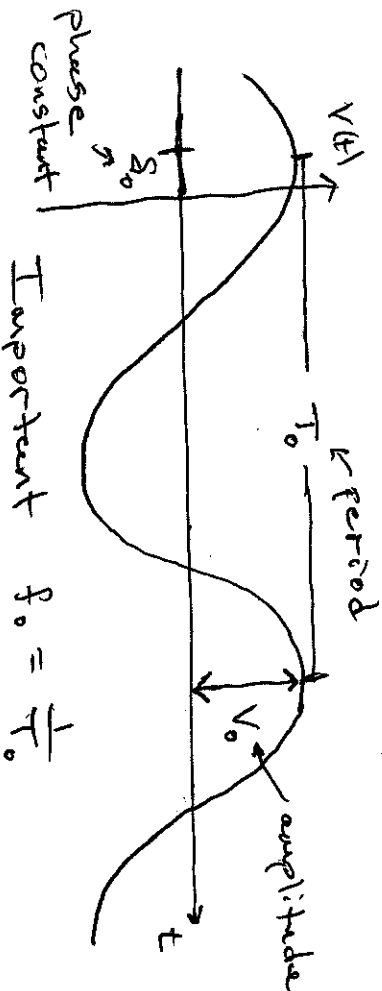
A physical oscillator has noise.

Remarkably all the important

sources of noise can be summarized in a single correction

$$V(t) = V_0 \cos(2\pi f_0 t + S_0 + \phi(t))$$

the time variation of the phase



A clock counts the oscillations

(say the peaks) of the oscillator

to measure how much time has

passed. Clearly  $\phi(t)$  can

mess this up, if it changes

the duration between the

peaks. One way to capture

this is to compare the physical

clock's total phase to the

ideal clock's phase, that is

Consider

$$\frac{\Phi(t)}{2\pi f_0} = t + \frac{S_0}{2\pi f_0} + \frac{\phi(t)}{2\pi f_0}$$

The  $S_0/2\pi f_0$  is a constant, we can always get rid of it by starting our clock at the right time. But,

$$x(t) = \frac{\phi(t)}{2\pi f_0} \quad \text{"relative phase"}$$

tells us how much our physical clock's time differs from a perfect

Dividing both sides by  $2\pi$  gives,

$$f = \frac{1}{2\pi} \frac{\Delta\Phi}{\Delta t} = f_0 + \frac{1}{2\pi} \frac{\Delta\phi}{\Delta t}$$

Then,

$$\Delta f = f - f_0 = \frac{1}{2\pi} \frac{\Delta\phi}{\Delta t}$$

So, we see that, <sup>the</sup> difference  $\Delta\phi$  also characterizes how our physical oscillator's frequency differs from the ideal oscillator's frequency  $f_0$ .

clock of frequency  $f_0$ , Pg/3

Before Spring break we also saw how to extract the frequency of a clock:

$$\begin{aligned} \frac{\Delta\Phi}{\Delta t} &= \frac{\Phi(t_2) - \Phi(t_1)}{t_2 - t_1} \\ &= \frac{(2\pi f_0 t_2 + S_0 + \phi(t_2)) - (2\pi f_0 t_1 + S_0 + \phi(t_1))}{t_2 - t_1} \\ &= \frac{2\pi f_0 (t_2 - t_1) + \phi(t_2) - \phi(t_1)}{t_2 - t_1} \end{aligned}$$

In practice there are several reasons that it is even better to work with

$$y = \frac{\Delta f}{f_0} \quad \text{"the fractional frequency difference"}$$

$$= 1 - \frac{f}{f_0} \quad \text{A pure \#}$$

$$= \frac{1}{2\pi f_0} \frac{\Delta\phi}{\Delta t} = \left( \frac{\Delta\phi}{2\pi f_0} \right) = \frac{\Delta x}{\Delta t}$$

The last expression,

$$y = \frac{\Delta x}{\Delta t}$$

will give us a very practical way to measure  $y$  in the lab.

More on that soon.

Why do we focus on  $y$ ? First and foremost it is because  $f$  is what we can measure. We take

There is another reason too: Why not just consider  $\Delta f = f - f_0$ ?

Suppose we had a pair of oscillators with

$$f = 11 \text{ Hz} \text{ and } f_0 = 10 \text{ Hz}$$

and a second pair with

$$f = 101 \text{ Hz} \text{ and } f_0 = 100 \text{ Hz}.$$

Both pairs have  $\Delta f = 1 \text{ Hz}$ , but

The best oscillator we can

own and treat it as ideal

and having exact frequency  $f_0$ .

Then we measure the number of oscillations of our non-ideal oscillator in an oscillation of our ideal oscillator. This is precisely

$$T_0 f = \frac{f}{f_0}.$$

if my clock's counter counts one extra cycle of the 1st case, it is off by about  $1/10$  of a second, while in the second case it is only off by about  $1/100$  of a second.

This difference is perfectly captured by  $y$ :

$$\text{Case 1: } y = \frac{\Delta f}{f_0} = \frac{1}{10} = 0.1$$

$$\text{Case 2: } y = \frac{\Delta f}{f_0} = \frac{1}{100} = 0.01.$$