

# Today's Outline

I Last Lecture

II The Equivalence Principle

III Gravitational Time Dilation

IV Grav. Time Dilation as Geometry

Doppler Effect:

$$\omega' = \omega \frac{\sqrt{1-v^2}}{1-v\cos\alpha'}$$

Check out relativistic beaming in text.

II The Equivalence Principle

Recall the parable of Galileo and the leaning tower of Pisa.

# Lecture 5

Jan 31<sup>st</sup>, 2012

I Last Lecture

A/5

- Discussed S.R. Kinematics & dynamics
- Touched on light

Light:

$$E = \hbar\omega = |\vec{p}|$$

write,

$$\vec{p} = \hbar\vec{k} \quad \omega / |\vec{k}| = \omega$$

then

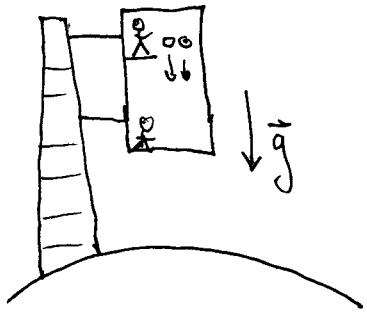
$$P^\alpha = (E, \vec{p}) = \hbar(\omega, \vec{k}) = \hbar k^\alpha$$

The importance of this story:

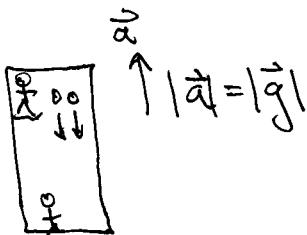
$m_I = m_G$

This lead to Einstein's "happiest thought", a free fall observer has no idea there's a gravitational field. Conversely, an

accelerated observer's experiments  
are identical to those of an  
observer immersed in a gravitational  
field pointing in the opposite direction:



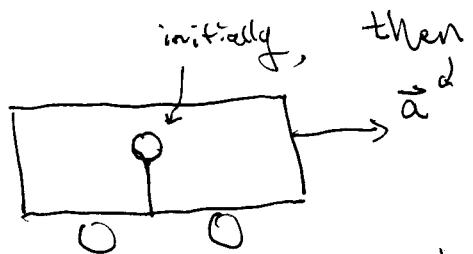
Same  
results  
as



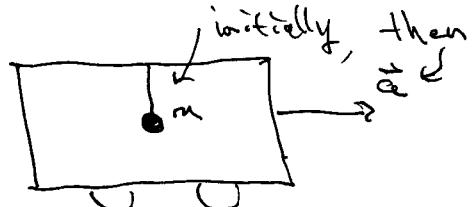
P2/5

Only an outside,  
inertial observer would  
say that these results  
happen for different  
reasons: gravitational  
attraction and rocket  
thruster respectively.

Example: Balloon in an accelerated  
train car.



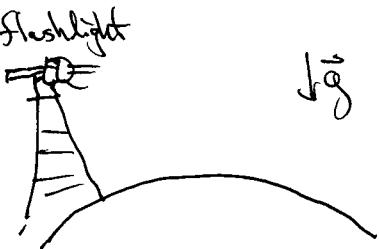
Which direction does the balloon  
deflect in? How about a pendulum?



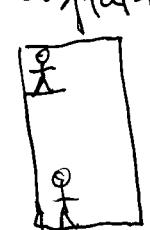
Easy to answer when you  
think of the gravitational  
analogs.

Works both ways:

Example: What happens to light  
rays in a grav. Field?



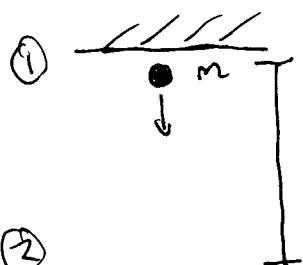
Same  
as



They're bent!

### III Gravitational Time Dilation

"Naive" approach: Drop a rock down an elevator shaft: Its kinetic energy is greater at the bottom,

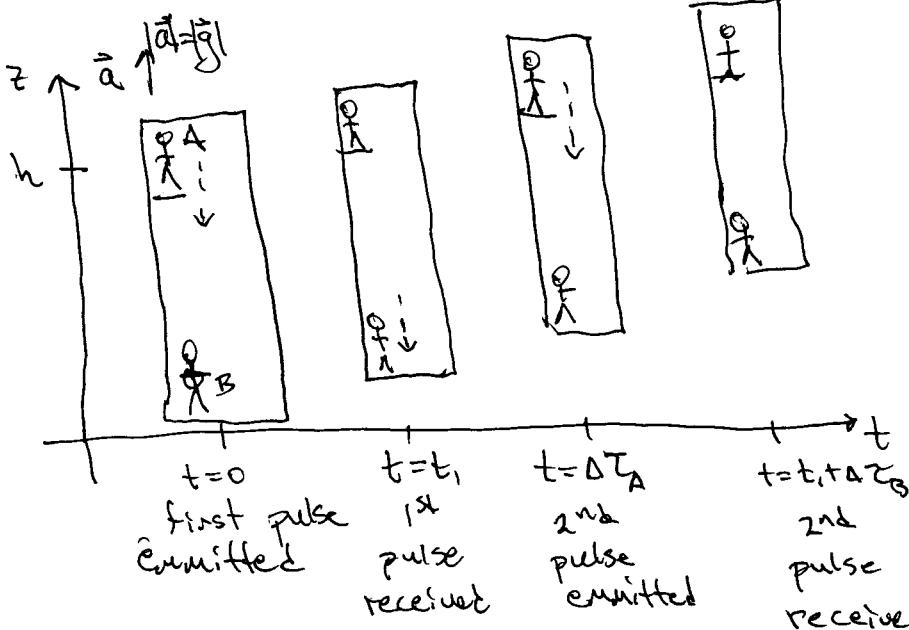


$$E_2 = E_1 + mgh$$

What if I drop a photon? Photons have  $E = \hbar\omega$ , so

$$\hbar\omega_2 = \hbar\omega_1 + \frac{\hbar\omega_1}{c^2} gh,$$

Quantitative approach: (Hartle)



where I've used  $m \rightarrow \infty$

$$\frac{E}{c^2} = \frac{\hbar\omega}{c^2}. \text{ So,}$$

$$\omega_2 = \omega_1 \left( 1 + \frac{gh}{c^2} \right)$$

The photon is blue-shifted by the fall. Confirmed experimentally by Pound & Rebka (1960). Tiny effect,

$$\frac{gh}{c^2} \approx \frac{10(20)}{9 \times 10^{12}} \approx 2 \times 10^{-15}$$

Neglect  $(\frac{v}{c})^2$ ,  $(\frac{gh}{c^2})^2$  and  $g(\frac{gh}{c^2})$  but keep  $\frac{v}{c}$  and  $\frac{gh}{c^2}$ .

$$z_B(t) = \frac{1}{2} gt^2$$

$$z_A(t) = h + \frac{1}{2} gt^2$$

Distance travelled by 1<sup>st</sup> pulse,

$$z_A(0) - z_B(t_1) = ct_1$$

by 2<sup>nd</sup> pulse,

$$z_A(\Delta T_B) - z_B(t_1 + \Delta T_B) = c(t_1 + \Delta T_B - \Delta T_A)$$

If  $\Delta\tau_A$  (hence  $\Delta\tau_B$ ) is small then,

$$h - \frac{1}{2}gt_1^2 = ct, \quad ①$$

$$\begin{aligned} h - \frac{1}{2}g(t_1 + \Delta\tau_B)^2 &\approx h - \frac{1}{2}gt_1^2 - gt_1\Delta\tau_B \\ &= c(t_1 + \Delta\tau_B - \Delta\tau_A) \quad ② \end{aligned}$$

And Eq ① - Eq ② is,

$$gt_1\Delta\tau_B = c(\Delta\tau_A - \Delta\tau_B)$$

$$\Rightarrow \Delta\tau_B = \frac{\Delta\tau_A}{(1 + gt_1/c)}$$

Reexpress in terms of gravitational potential (per unit mass, e.g.  $\Phi = -\frac{GM}{r}$   
not  $V = -\frac{GMm}{r}$  and  $[\Phi] = \frac{\text{Energy}}{\text{Mass}}$ )  $\Phi$ ,

$$\Phi_A - \Phi_B = gh$$

and rates of reception  $\text{rate}_A = \frac{1}{\Delta\tau_A}$ ,  
 $\text{rate}_B = \frac{1}{\Delta\tau_B}$ , then,

$$\text{rate}_B = \text{rate}_A \left(1 + \frac{gh}{c^2}\right) = \text{rate}_A \left(1 + \frac{\Phi_A - \Phi_B}{c^2}\right)$$

But from Eq. ①,

P1/5

$$\frac{1}{2}gt_1^2 + ct_1 - h = 0$$

$$\begin{aligned} \Rightarrow t_1 &= \frac{-c \pm \sqrt{c^2 + 2gh}}{g} \\ &= -\frac{c}{g} \pm \frac{c}{g} \left(1 + \frac{gh}{c^2} + \dots\right) \\ &= \frac{h}{c} + \dots \end{aligned}$$

$$\text{Thus, } \Delta\tau_B = \Delta\tau_A \left(1 - \frac{gh}{c^2}\right)$$

Can summarize as:

"Slouching clocks run slow"

#### IV Gravitational Time

Dilation as Geometry

We begin our program in earnest: I hand you a metric, say a little about it and then you

explore it.

Our First curved Spacetime metric:

$$ds^2 = -\left(1 + \frac{2\Phi(x^i)}{c^2}\right)(cdt)^2 + \left(1 - \frac{2\Phi(x^i)}{c^2}\right)(dx^2 + dy^2 + dz^2),$$

here  $\Phi$  is the gravitational potential in the Newtonian sense and is a

So this is the static weak field metric.

function of position alone, PS/5

$$\underline{\Phi} = \underline{\Phi}(x^i) \quad i=1, 2, 3$$

we call this a static potential. This fine element is only valid for small curvatures and weak sources, which is known as a weak field.