

# Math Methods

## Day 13

### Last Time

- We studied the limit  $N \rightarrow \infty$  of our model and found  $\frac{p}{N+1} = \frac{x}{L}$ , as well as

$$y(x,t) = A \sin(n\pi \frac{x}{L}) \cos(\omega_n t + \phi_n),$$

with

$$\omega_n = \frac{n\pi}{L} \sqrt{\frac{T}{\mu}}$$

- We derived the wave equation

$$\frac{\partial^2 y}{\partial t^2} = v^2 \frac{\partial^2 y}{\partial x^2}$$

for waves on a taut string and showed  $v = \sqrt{T/\mu}$ .

### Normal Modes and the Wave Equation

**Normal Modes:** The definition of a normal mode is that each position along the string oscillates at the same frequency. We can encode this definition in a good guess for the solution of the wave equation:

$$y(x,t) = f(x) \cdot \cos(\omega t + \phi).$$

Here  $f(x)$  is an "amplitude function" that we don't know how to anticipate and which we will try to adjust to solve the wave equation.

Using this guess,

$$\begin{aligned} \frac{\partial^2 y}{\partial t^2} &= -\omega^2 f(x) \cos(\omega t + \phi) \\ \frac{\partial^2 y}{\partial x^2} &= \frac{d^2 f}{dx^2} \cos(\omega t + \phi). \end{aligned}$$

Putting these expressions into the wave equation gives:

$$\begin{aligned} -\omega^2 f(x) \cos(\omega t + \phi) &= v^2 \frac{d^2 f}{dx^2} \cos(\omega t + \phi) \\ \implies \frac{d^2 f}{dx^2} &= -\left(\frac{\omega}{v}\right)^2 f \end{aligned}$$

So,

$$f(x) = A \sin\left(\frac{\omega}{v}x\right) + B \cos\left(\frac{\omega}{v}x\right),$$

and

$$y(x,t) = \left[ A \sin\left(\frac{\omega}{v}x\right) + B \cos\left(\frac{\omega}{v}x\right) \right] \cos(\omega t + \phi).$$

### Today

- I. Last Time
- II. Normal Modes & the Wave Equation
- III. General Solutions & Fourier Series

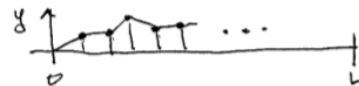


Figure 1: Our discrete model of transverse oscillations.

$\mu$  = linear mass density



Figure 2: A taut string of length L.

Simple Harmonic Oscillator Eq. (SHO)

But: we have boundary conditions! Recall that we said that the string was tacked down at both ends, then

$$(1) \quad y(0, t) = 0; \quad (2) \quad y(L, t) = 0.$$

From (1), and since this boundary condition is supposed to hold for all times  $t$ ,

$$B \cos(\omega t + \phi) = 0 \implies B = 0.$$

So,

$$y(x, t) = A \sin\left(\frac{\omega}{v}x\right) \cos(\omega t + \phi).$$

Then (2) implies

$$y(L, t) = A \sin\left(\frac{\omega}{v}L\right) \cos(\omega t + \phi) = 0.$$

We don't want  $A = 0$ , so require

$$\sin\left(\frac{\omega}{v}L\right) = 0 \implies \frac{\omega}{v}L = n\pi \quad (n = 1, 2, 3, \dots),$$

or

$$\omega_n = \frac{n\pi}{L}v = \frac{n\pi}{L}\sqrt{\frac{T}{\mu}}.$$

This is in perfect agreement with our limit of the coupled oscillators. The normal modes are

$$y_n(x, t) = A \sin\left(\frac{n\pi}{L}x\right) \cos(\omega_n t + \phi)$$

$$\omega_n = \frac{n\pi v}{L}; \quad n = 1, 2, 3, \dots,$$

these are standing waves!

Note:

$$\omega_n = \frac{n\pi v}{L} \implies v = \frac{L\omega_n}{n\pi} = \frac{2L}{n} \frac{\omega_n}{2\pi} = \lambda f.$$

BUT: Normal modes are not the most general motion of a string. Two nice examples that clearly don't give standing waves are:

- A plucked guitar string
- A hammered piano string

### General Solutions & Fourier Series

The General Solution: The two pictures at right correspond to interesting initial conditions:

$$F(x) = y(x, 0) \quad (\text{given initial shape})$$

and

$$G(x) = \dot{y}(x, 0) = \frac{\partial y}{\partial t}(x, 0) \quad (\text{given initial velocity}).$$

$n^{\text{th}}$  normal mode frequency

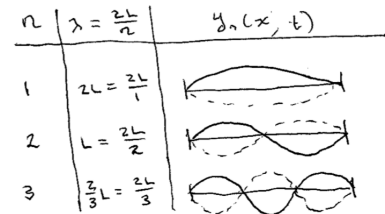


Figure 3: The first three normal modes.

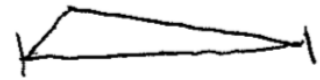


Figure 4: A plucked string has a rich initial shape, but no initial velocity.

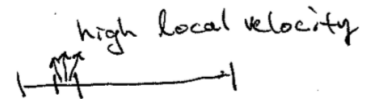


Figure 5: A hammered string has a rich initial velocity, but no initial deformation.

Claim: The general solution is (as for coupled oscillators) a linear combination of the normal modes. Namely,

$$y(x, t) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi}{L}x\right) \cos\left(\frac{n\pi v}{L}t + \phi_n\right).$$

Our two special initial conditions require either

$$F(x) = y(x, 0) \implies F(x) = \sum_{n=1}^{\infty} (A_n \cos \phi_n) \sin\left(\frac{n\pi}{L}x\right),$$

or

$$G(x) = \dot{y}(x, 0) \implies G(x) = \sum_{n=1}^{\infty} \left(-A_n \frac{n\pi v}{L} \sin \phi_n\right) \sin\left(\frac{n\pi}{L}x\right).$$

We have to pick  $A_n$  and  $\phi_n$  so as to meet these two conditions. Can we do this? How?

We go to math. To do so, let's recast this as a question directly to mathematics: given  $f(x)$ , find constants  $b_n$  such that

$$f(x) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi}{L}x\right) \quad \text{on } (0 \leq x \leq L).$$

This expression is called the Fourier series expansion for  $f(x)$ .