

Outline

I Lorentz transformations:
velocity addition & the Spacetime
invariant

II Four - Vectors

III Energy and Momentum
in Relativity

Modern
Day 4

I. Recall the Lorentz
transformations:

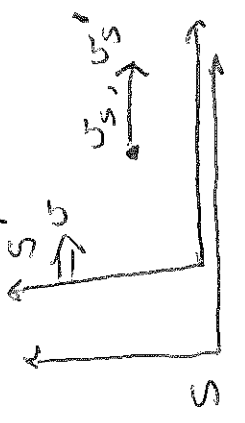
$$\begin{aligned}
 x' &= \gamma (x - \frac{v}{c} (ct)) \\
 ct' &= \gamma (ct - \frac{v}{c} x)
 \end{aligned}$$

Of course,

$$\begin{aligned}
 y &= y' \\
 z &= z'
 \end{aligned}$$

This half page intentionally
blank

Velocity Addition



$$\left. \begin{aligned} \Delta x = x_2 - x_1 \\ \Delta t = t_2 - t_1 \end{aligned} \right\} \Rightarrow v_s = \frac{\Delta x}{\Delta t}$$

In S' $\left. \begin{aligned} \Delta x' = x'_2 - x'_1 = \gamma(\Delta x - v\Delta t) \\ \Delta t' = t'_2 - t'_1 = \gamma(\Delta t - \frac{v}{c^2}\Delta x) \end{aligned} \right\} v'_s = \frac{\Delta x'}{\Delta t'}$

$$v'_s = \frac{\gamma(\Delta x - v\Delta t)}{\gamma(\Delta t - \frac{v}{c^2}\Delta x)} = \frac{\frac{\Delta x}{\Delta t} - v}{1 - \frac{v}{c^2}\frac{\Delta x}{\Delta t}}$$

We have $\vec{a} = (1, 0)$ and

$$\vec{a}' = (\cos\theta, \sin\theta)$$

Geometrically it is clear that

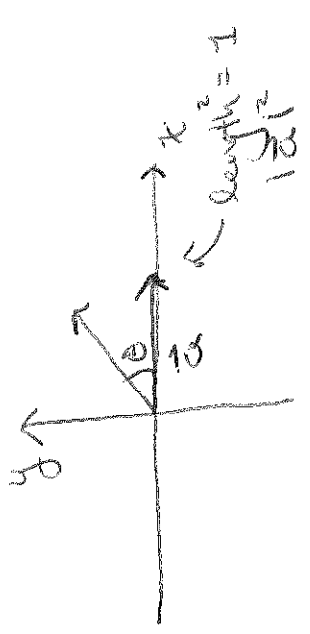
$$|\vec{a}'|^2 = 1 \Rightarrow \cos^2\theta + \sin^2\theta = 1 \checkmark$$

Rotations leave the length of an (spatial) interval invariant. What about Lorentz transforms? Consider

$$(\Delta x')^2 - c^2 \Delta t'^2 = \gamma^2 [\Delta x^2 - 2v\Delta x\Delta t + v^2\Delta t^2 - c^2\Delta t^2]$$

Spacetime Invariant
A very useful way to PZ/4 think about transformations is to ask what, if anything, they leave invariant.

Consider rotations



$$\begin{aligned} & -c^2 (\Delta t^2 - \frac{2v}{c^2} \Delta x \Delta t + \frac{v^2}{c^2} \Delta x^2) \\ & = \gamma^2 \left[(1 - \frac{v^2}{c^2}) \Delta x^2 - (c^2 - v^2) \Delta t^2 \right] \\ & = \gamma^2 \left[\frac{1}{\gamma^2} \Delta x^2 - (1 - \frac{v^2}{c^2}) c^2 \Delta t^2 \right] \\ & = \Delta x^2 - c^2 \Delta t^2 \end{aligned}$$

The interval

$$\Delta x^2 - c^2 \Delta t^2$$

is invariant, i.e. the same in any frame of reference.

Tibber Notation:

$x^0 \equiv ct, x^1 \equiv x, x^2 \equiv y, x^3 \equiv z$

$(x^0)' = \gamma(x^0 - \beta x^1), \beta \equiv \frac{v}{c}$

$(x^1)' = \gamma(x^1 - \beta x^0)$

$(x^2)' = x^2; (x^3)' = x^3$

Position-time 4-vector:

$x^\mu = (x^0, x^1, x^2, x^3) = (ct, x, y, z)$

Let's build up another example.

Proper time: Two times: "ordinary"

("coordinate") time, shown on

lab wall clock - t .



$\begin{matrix} \text{lab} \\ \rightarrow \\ \text{wrist watch time} \end{matrix}$

(wrist watch time is) "proper" - shown

on a particle's own watch: τ .

II Four-Vectors: A "4-vector" P3/14

is any collection of 4 #'s

$a^\mu = (a^0, a^1, a^2, a^3)$ that

transforms in the same way that x^μ does when you go from

S to S'

$(a^0)' = \gamma(a^0 - \beta a^1)$

$(a^1)' = \gamma(a^1 - \beta a^0)$

$(a^2)' = a^2; (a^3)' = a^3$.

(Note: Proper time is same for all observers; t is not.)

$dt = \gamma d\tau$

Proper Velocity: "Ordinary"

velocity: $U = \frac{\Delta x}{\Delta t}$ (both lab system)

"proper" velocity:

$\frac{\Delta x}{\Delta \tau}$ (lab system)

$v = \frac{\Delta x}{\Delta t}$ (moving system)

Notice:

$$v = \frac{\Delta x}{\Delta t} = \frac{\Delta x}{\gamma \Delta \tau} = \frac{1}{\gamma} v$$

Proper velocity 4-vector

$$\eta^\mu = \frac{dx^\mu}{d\tau}$$

$$\eta^0 = \frac{dx^0}{d\tau} = \frac{c dt}{d\tau} = \gamma c$$

$$\eta^1 = \frac{dx^1}{d\tau} = \frac{dx}{d\tau} = \gamma v_x$$

$$\eta^2 = \frac{dx^2}{d\tau} = \gamma v_y$$

p4/4

$$\text{So, } \eta^\mu = (c\gamma, \gamma v_x, \gamma v_y, \gamma v_z)$$

$$= \gamma (c, v_x, v_y, v_z)$$

Note that

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \Rightarrow \gamma^2 = \frac{1}{1 - \frac{v^2}{c^2}}$$

$$\Rightarrow \gamma^2 \left(1 - \frac{v^2}{c^2}\right) = 1 \Rightarrow \gamma^2 = 1 + \frac{v^2}{c^2}$$

$$\Rightarrow \gamma = \sqrt{1 + \frac{v^2}{c^2}}$$

III Momentum: classically: mass * velocity

$$\vec{p} = m \vec{v}, \text{ or } \vec{p} = m \vec{\eta}?$$

↳ This def. would \Rightarrow that conservation of momentum is not even a possible law \leftarrow if in S, not true in S'.

Define: Relativistic momentum:

$$\vec{p} = m \vec{\eta} = \gamma m \vec{v} = \frac{m \vec{v}}{\sqrt{1 - \frac{v^2}{c^2}}}$$

This means that \vec{p} is the spatial

part of a 4-vector:

$$p^\mu = m \eta^\mu$$

Question: What's the zeroth ("temporal") component?

$$p^0 = m \eta^0 = \gamma m c$$

Define: Relativistic energy

$$E = \gamma m c^2$$

$$\text{or } E = \frac{m c^2}{\sqrt{1 - \frac{v^2}{c^2}}}$$