

Outline

I Last time

II Michelson-Morley Experiment

III Wave motion

Modern

Day 8

I. Conservation of E and P

→ hold in relativity
provided you use their
relativistic definitions.

In doing calculations:
Suggestion 2: Work with E and \vec{P} .
Suggestion 1: Use m , not μ .

$$E^2 - \vec{P}^2 c^2 = m^2 c^4$$

to get E given \vec{P} and vice versa.

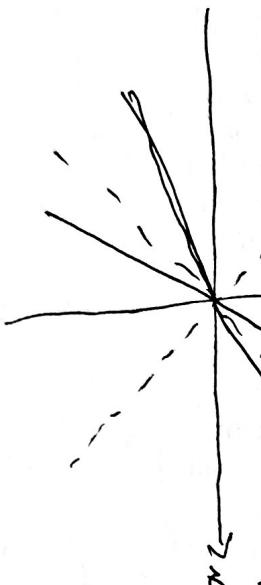
Suggestion 2: To get \vec{U} give E and \vec{P}
use,

$$\vec{U} = \frac{\vec{P} c^2}{E}$$

Can depict a boosted frame using
at t , t' slope = $\frac{c}{U} = \beta$

have treated light as an
object, much like a particle.

Part of the discovery of
light's character came through
its wave properties, to
which we turn now.



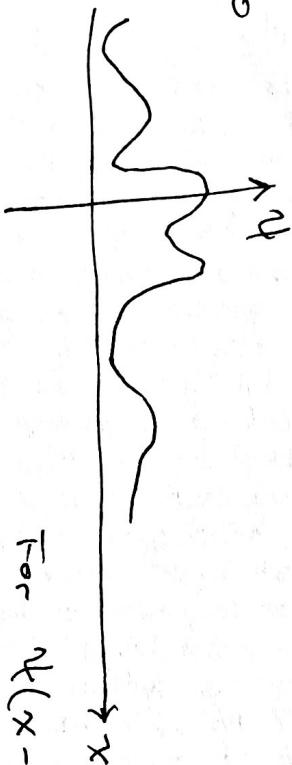
II Einstein's 2nd Postulate
highlights the special nature
of light. But, so far we

Part of the discovery of
light's character came through
its wave properties, to
which we turn now.

The waves most familiar from our surroundings are the motion of energy and momentum through a medium; e.g. water, air, in wood (marimba), through rock (earthquake).

If light is an electromagnetic wave, what is the medium it is waving? This is the question that Michelson and Morley

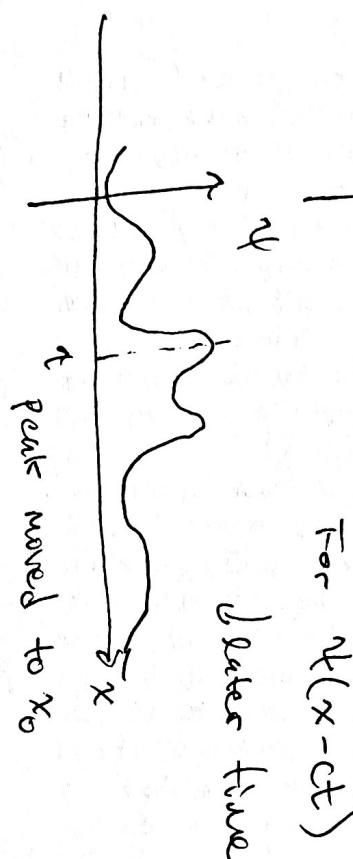
$$t=0$$



$$\psi(x+ct) \text{ solns are left-moving waves.}$$

For $\psi(x-ct)$

$$t_0 > 0$$



General property: superposition
if $\psi_1(x, t)$ & $\psi_2(x, t)$ solns
the $\psi_1(x, t) + \psi_2(x, t)$ is too.

didn't ask. But their experiment P2/3 hinted at the answer.

III Wave motion

The wave equation (in 1D)

$$\frac{\partial^2 \psi(x, t)}{\partial t^2} = c^2 \frac{\partial^2 \psi(x, t)}{\partial x^2}$$

c is a constant \rightarrow speed of wave.
Any differentiable function of the form $\psi(x-ct)$ or $\psi(x+ct)$ satisfies the wave equation.

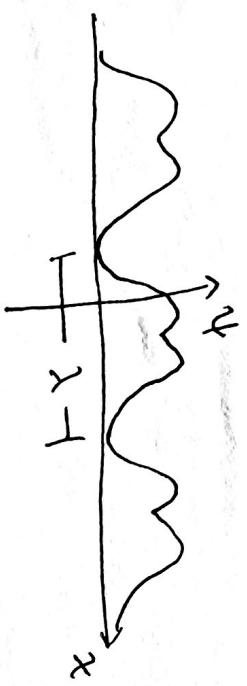
$$\text{We see, } \frac{x_0}{t_0} = c \text{ (positive) speed of wave.}$$

$$\psi(x+x_0, t_0) = \psi(x, 0)$$

$$\Rightarrow \psi(x+x_0 - ct_0) = \psi(x) \Rightarrow x - ct_0 = 0$$

A (useful) subset of solns are the periodic waves:

$$(Eq 2) \quad \psi(x+\lambda, t) = \psi(x, t)$$



λ - wave length, but then

$$f = \frac{c}{\lambda} - \text{call it frequency}$$

$$\frac{1}{f} = T = \frac{\lambda}{c} \quad \text{Period.}$$

$$\text{Eq: } \lambda \Rightarrow x + \lambda - ct = x - ct = x - c(t - \tau)$$

$$\Rightarrow \psi(x, t - \tau) = \psi(x, t) \quad \begin{matrix} \text{Periodic} \\ \text{in time} \end{matrix}$$

Usual form

$$A \cos(kx \pm kt + \phi) = A \cos(k(cx + \lambda) \pm kt + \phi)$$

$$\Rightarrow k = \frac{2\pi}{\lambda} \quad \text{and} \quad k_c = \frac{2\pi}{T} = \omega \quad \text{"angular freq!"}$$

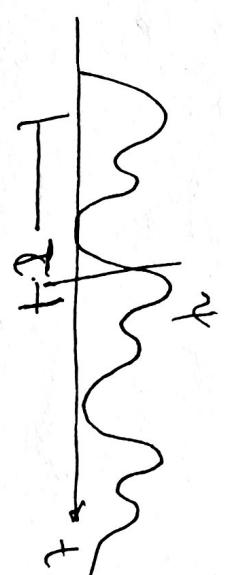
$$= A \cos(kx \pm \omega t + \phi)$$

Harmonic waves are useful:

1. sin/cos are easy to manipulate

2. Any periodic wave $\psi =$ sum (possibly ∞) of harmonic waves (Fourier Series)

3. Any wave over finite range of x but can be written as sum of harmonic waves



Further subset: harmonic waves
(sinusoidal solns)

$$\psi(x, t) = A \sin(x \pm ct), B \cos(x \pm ct)$$

Of course, sin and cos are just offsets of one another, we write

$$A \cos(\frac{\pi}{2}[x \pm ct] + \phi)$$

Amplitude wave # Phase or phase shift

Even easier to use than sin's and cos's

$$\boxed{e^{\pm i\theta} = \cos\theta \pm i\sin\theta} \quad \begin{matrix} \text{Euler's} \\ \text{relation} \end{matrix}$$

$$\Rightarrow \cos\theta = \frac{e^{i\theta} + e^{-i\theta}}{2}, \sin\theta = \frac{e^{i\theta} - e^{-i\theta}}{2i}$$

This will be our theme

next time.