

Today

I. Last Time

II. Office Hour Change Tuesdays, HW Folders, & Lab

III. More on Spacetime (Minkowski) Diagrams

IV. Galilean and Lorentz Transformations

I. We finished deriving three of the four fundamental consequence of Einstein's two postulates.

1. Relativity of simultaneity

2. Time dilation: $\Delta t = \gamma \Delta t'$

3. Length (or Lorentz) contraction: $\Delta L = \frac{1}{\gamma} \Delta L'$

4. You are deriving Einstein velocity addition on the homework.

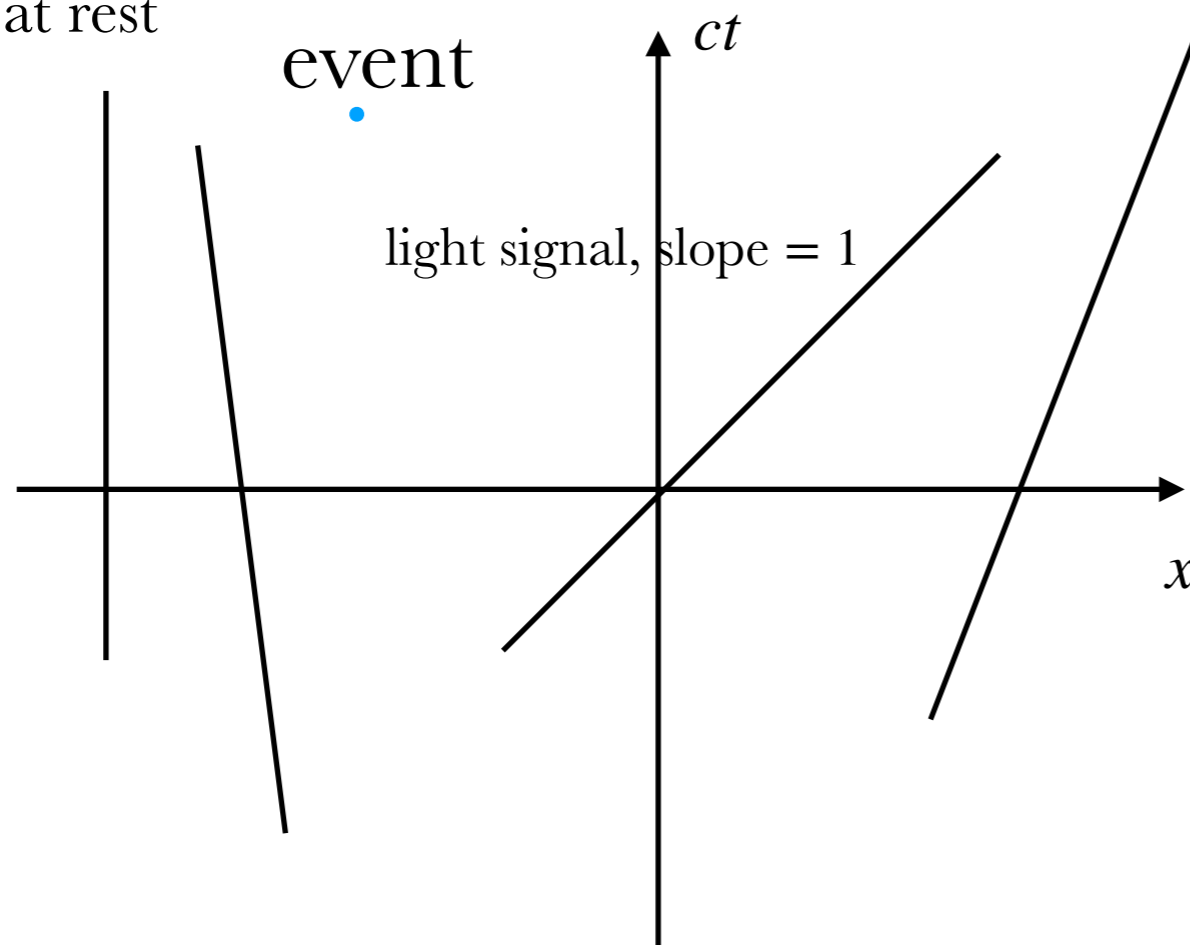
2. and 3. are starting to hint that maybe space and time fit together.

We won't derive the 4th consequence here, you're doing it on your homework. This is the Einstein velocity addition formula.

III. More on Spacetime (Minkowski) Diagram

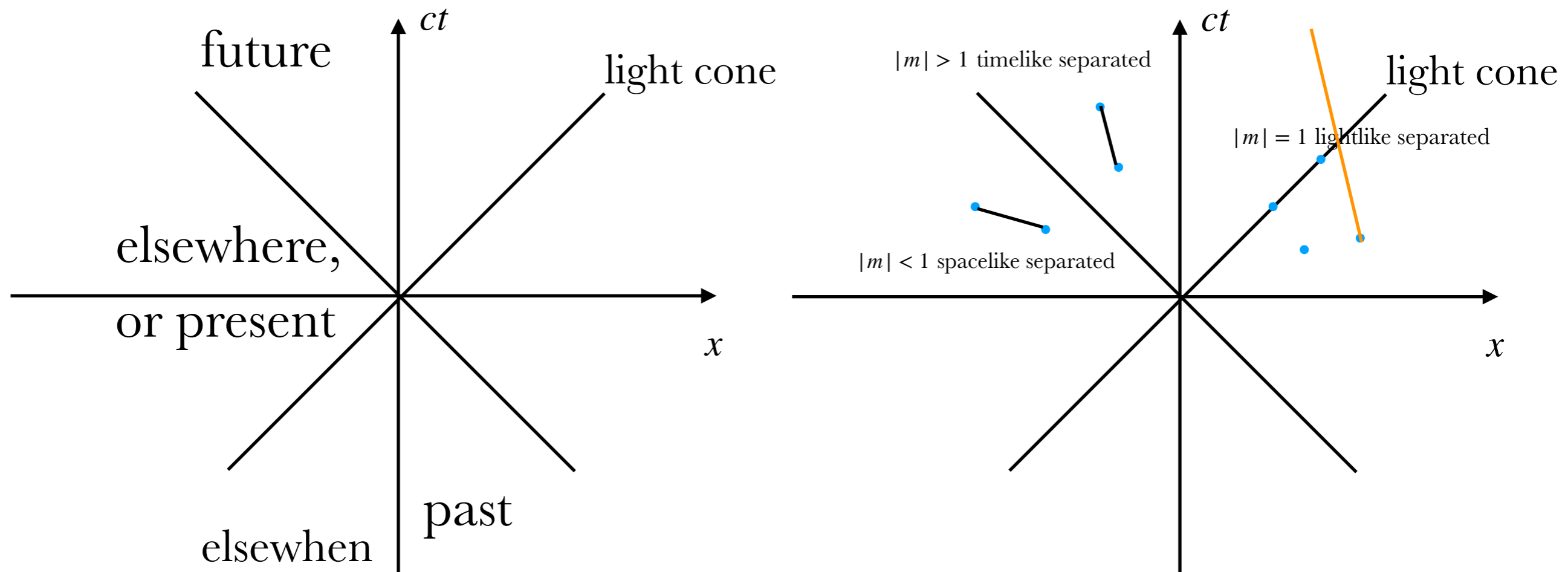
$$\text{slope} = \frac{\text{rise}}{\text{run}} = \frac{c\Delta t}{\Delta x} = \frac{c}{\Delta x/\Delta t} = \frac{c}{v}$$

Particle at rest



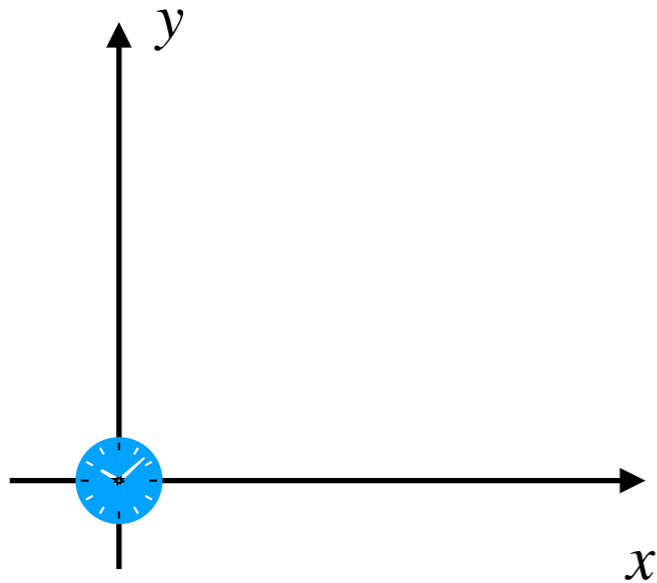
$$\frac{v}{c} = \frac{1}{\text{slope}}$$

III. Spacetime (Minkowski) Diagram

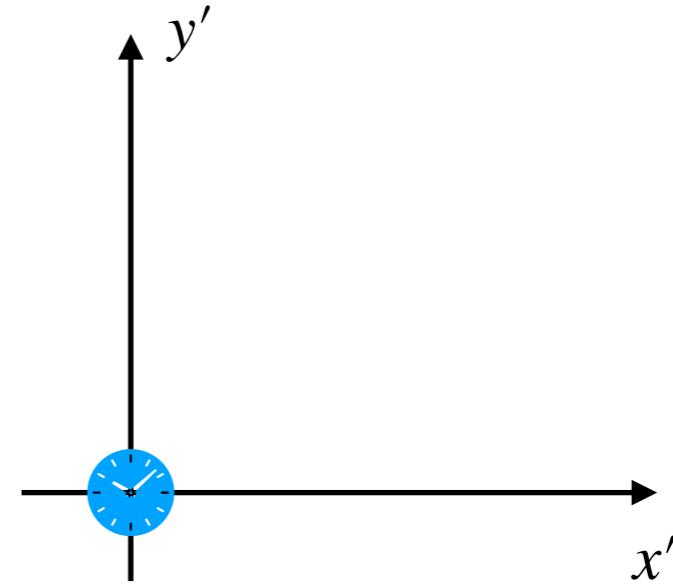


All of the diagrams that I've drawn so far are making reference to just one reference frame. Ultimately we want the dictionary that connects different frames of reference. This dictionary is called the Lorentz transformation.

IV. Galilean and Lorentz Transformations



Frame S (“rest”,
“ground”, “lab”)



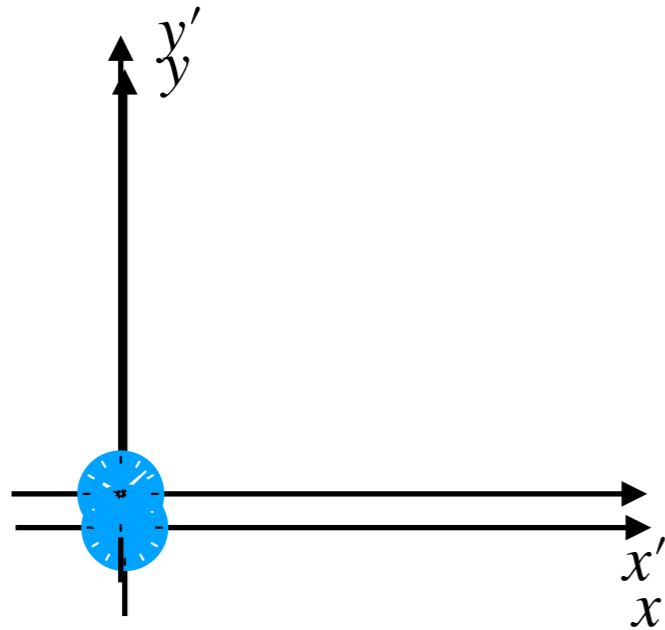
Frame S' (“moving”
frame, “train”, “rocket”)

Event: particular spacetime locations (x, y, z) and t (all in S).

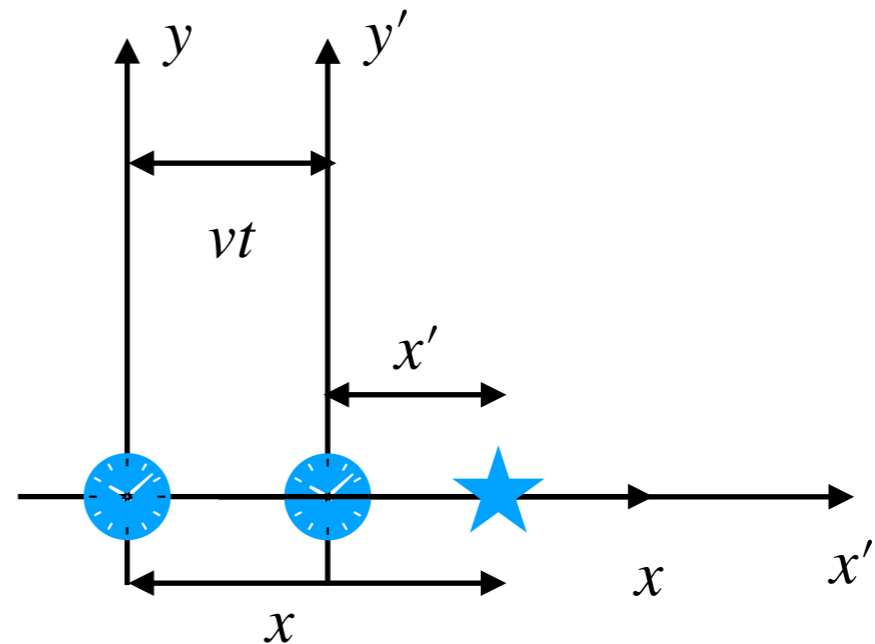
Question: What are the coordinates (x', y', z', t') of the SAME EVENT in S' ?

Let's assume that frame S' moves at speed v (const.) relative to S , along the x -axis.

IV. Galilean and Lorentz Transformations



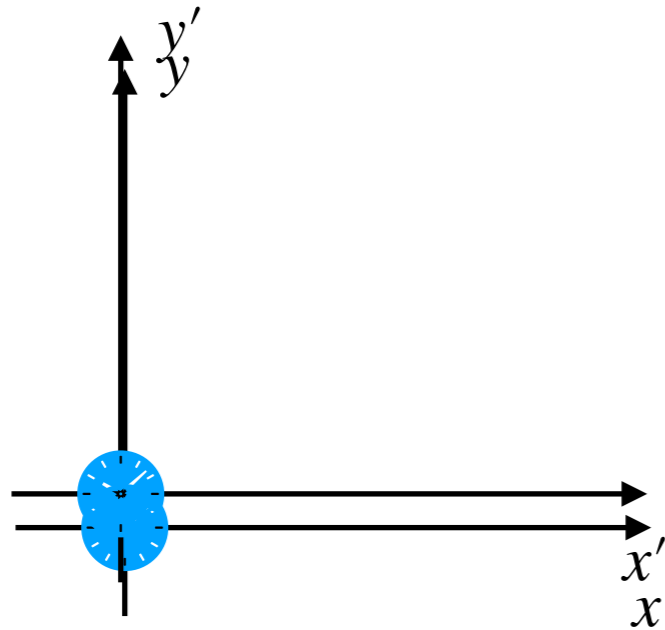
Initially



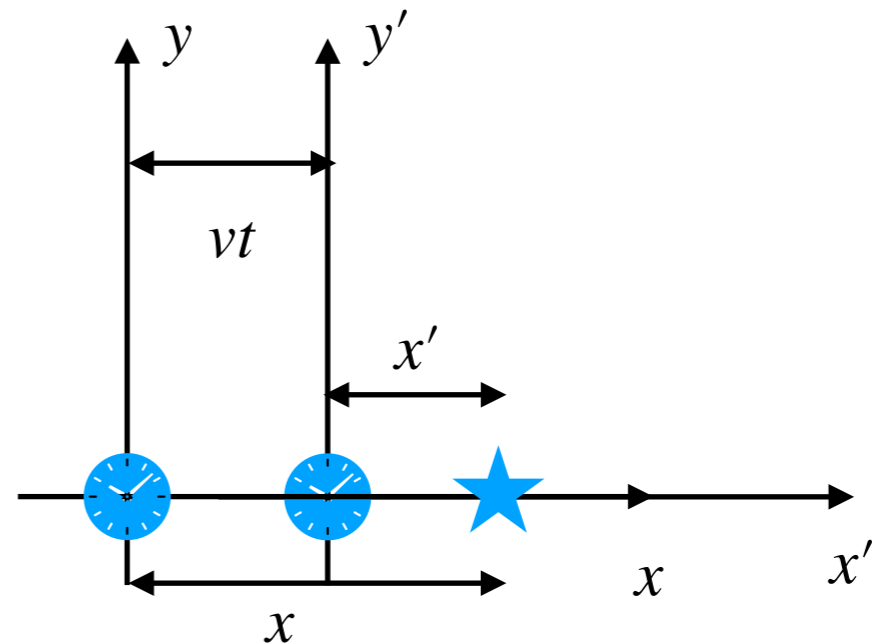
Frame S' (“moving”
frame, “train”, “rocket”)

Let's assume that frame S' moves at speed v (const.) relative to S , along the x -axis. At the instant the two origins coincide, we set the two master clocks to zero. After time t , the origin of S' has moved a distance vt .

IV. Galilean and Lorentz Transformations



Initially



Frame S' (“moving”
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For Galilean transformations we have

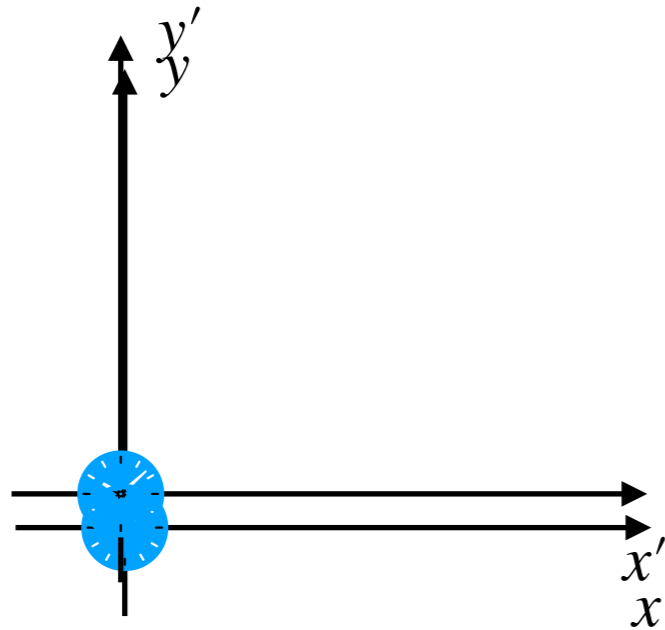
$$x' = x - vt$$

$$y' = y$$

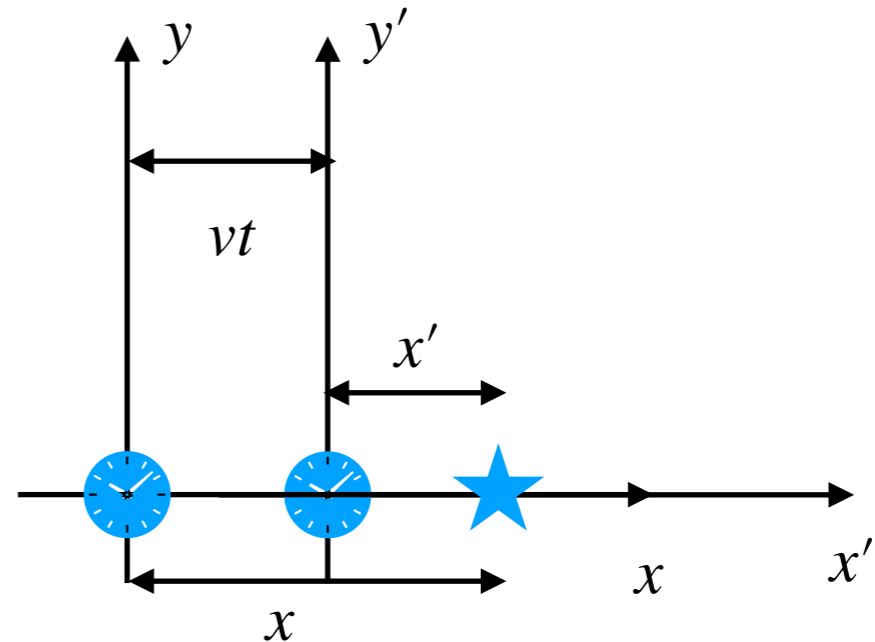
$$z' = z$$

$$t' = t$$

The Lorentz Transformations



Initially



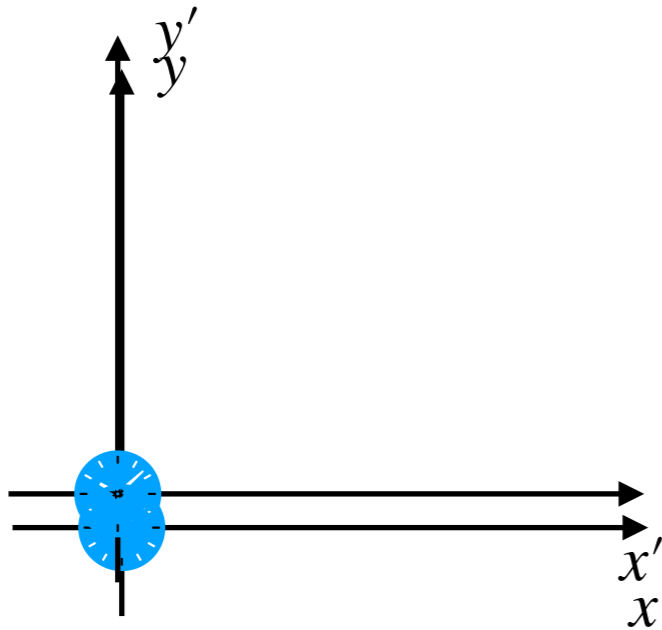
Frame S' (“moving”
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Take perspective of S . The Galilean “ x ” failed to account for length contraction. Moving objects are shorter than in their rest frame, so

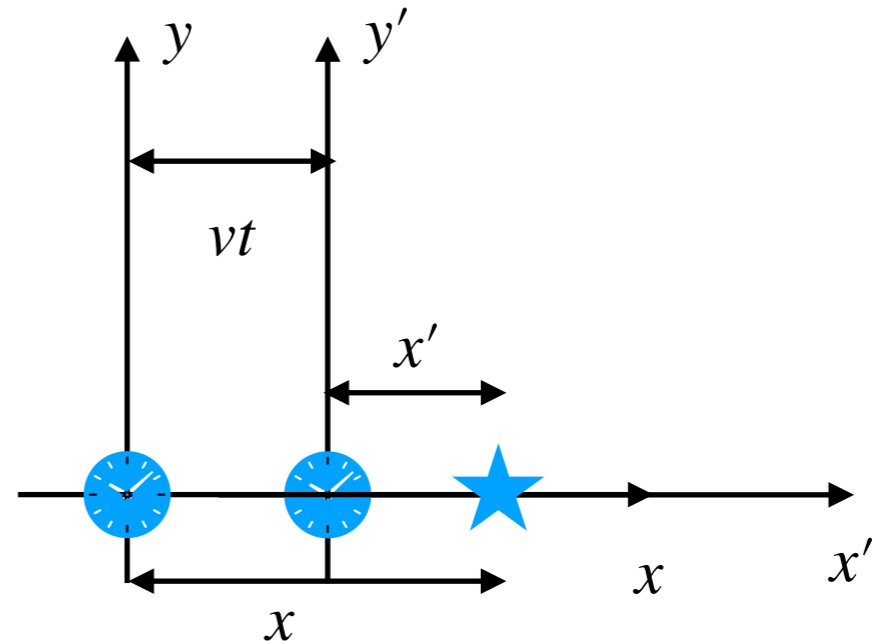
$$x' = \gamma(x - vt)$$

$$y' = y, z' = z. \text{ Not done yet(!), need time } t' = \gamma\left(t - \frac{v}{c^2}x\right)$$

The Lorentz Transformations



Initially



Frame S' (“moving”
frame, “train”, “rocket”)

Let's derive this: $t' = \gamma(t - \frac{v}{c^2}x)$. Inverse transformation of

position: $x = \gamma(x' + vt')$ $= \gamma[\gamma(x - vt) + vt'] = \gamma^2 x - \gamma^2 vt + \gamma vt'$.

$$\gamma vt' = (1 - \gamma^2)x + \gamma^2 vt$$

$$t' = \frac{(1 - \gamma^2)}{\gamma v} x + \gamma t$$

The Lorentz Transformations

Let's derive this: $t' = \gamma(t - \frac{v}{c^2}x)$. Inverse transformation of

position: $x = \gamma(x' + vt')$ $= \gamma[\gamma(x - vt) + vt'] = \gamma^2x - \gamma^2vt + \gamma vt'$.

$$\gamma vt' = (1 - \gamma^2)x + \gamma^2vt$$

$$t' = \frac{(1 - \gamma^2)}{\gamma v}x + \gamma t$$

Focusing on the term in front of x

$$\frac{1}{v} \frac{1 - \gamma^2}{\gamma} = \frac{\gamma}{v} \frac{1 - \gamma^2}{\gamma^2} = \frac{\gamma}{v} \left(\frac{1}{\gamma^2} - 1 \right) = \frac{\gamma}{v} \left(1 - \frac{v^2}{c^2} - 1 \right) = -\frac{\gamma v}{c^2}$$

Plugging this in gives

$$t' = \gamma t - \frac{\gamma v}{c^2}x = \gamma \left(t - \frac{v}{c^2}x \right).$$

In summary, we have two results

$$x' = \gamma(x - vt) \quad \text{and} \quad t' = \gamma \left(t - \frac{v}{c^2}x \right)$$

Time meets	Space
2-2:45	2:40-3:20
Spencer	Bobby
Woochan	Jade
Skylar	Rose
Ollie	Thomas
Talia	Grace
Antonio	Quincy