<u>Today</u>

- I. Last Time
- II. Interference of waves with a phase difference
- III. Complex index of refraction
- IV. Diffusion Equation
 - I. Yanpei Deng this week due to exam (will help in the lab. She's available MTuW from 7-8pm in Brody lab).
 This week: Antu Antu will be providing homework support. Hours are: Tu 8-9pm, Th10:30-11:30am, Th 8-9pm.

Today's office hours 5:30-7pm. Thursday office hours will now run 5-6:30pm. I. Last time we started to discuss the intensity of a superposition of harmonic waves. We ended right in the midst of that calculation and part of today will be devoted to completing it.

Discussed attenuation in a lot of detail. Made a connection with our first M&M lab where we had a fixed number of decays for each trial. In the case of waves we lose a fixed percentage of the amplitude for every unit of distance that we travel into the medium:

 $A(x) = A(0)e^{-\alpha x}$, (attenuation equation).

Of course, it is only approximately true in some specific materials that we lose the same amplitude at every step. We describe this attenuation with the "characteristic length" or "skin depth", which is $\ell = 1/\alpha$.

We noticed that it is ω that it constant across a boundary, not λ .

II. What is the intensity of a superposition of harmonic waves with same E_0 , k, ω , but different ϕ ?

First, the wave is as follows

$$E(x,t) = E_0 \cos(kx - \omega t) + E_0 \cos(kx - \omega t + \phi).$$

The intensity is

$$I = \epsilon_0 c \overline{E^2} = \epsilon_0 c E_0^2 \overline{[\cos(kx - \omega t) + \cos(kx - \omega t + \phi)]^2}$$

$$I = \epsilon_0 c E_0^2 \overline{[\cos^2(kx - \omega t) + \cos^2(kx - \omega t + \phi) + 2\cos(kx - \omega t)\cos(kx - \omega t + \phi)]}$$

$$I = \epsilon_0 c E_0^2 [\frac{1}{2} + \frac{1}{2} + 2\overline{\cos(kx - \omega t)\cos(kx - \omega t + \phi)}]$$

Going from here we have to figure out what to do with this last term: recall $\cos A \cos B = \frac{1}{2} \cos(A+B) + \frac{1}{2} \cos(A-B)$. Then $I = \epsilon_0 c E_0^2 [\frac{1}{2} + \frac{1}{2} + 2\left[\frac{1}{2} \cos(-\phi) + \frac{1}{2} \cos(2kx - 2\omega t + \phi)\right]]$ $I = \epsilon_0 c E_0^2 [1 + \cos \phi]$

For $\phi = 0$ we get twice the amplitude and four times the intensity!

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 $I = \epsilon_0 c E_0^2 [1 + \cos \phi]$
For $\phi = 0$ we get twice the amplitude and four times the intensity!
The case $\phi = \pi$ gives completely destructive interference.

III. Let's return to our wave that experiences attenuation, A(x). Let's also assume that it's a harmonic wave

$$\begin{split} \psi(x,t) &= A(x)\cos(kx - \omega t) = A_0 e^{-\alpha x}\cos(kx - \omega t) \\ & \Rightarrow \tilde{\psi}(x,t) = A_0 e^{-\alpha x} e^{i(kx - \omega t)} = A_0 e^{-\alpha x + i(kx - \omega t)} = A_0 e^{i(i\alpha x + kx - \omega t)} = A_0 e^{i([k+i\alpha]x - \omega t)} \\ & \text{This is just like the harmonic waves we started with if we let} \\ & \tilde{k} = k + i\alpha. \text{ Then we have a super compact wave: } \tilde{\psi}(x,t) = A_0 e^{i(\tilde{k}x - \omega t)}. \end{split}$$

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Recall that the wave speed is given by:

 $c = \lambda f = \frac{2\pi}{k} \frac{\omega}{2\pi} = \frac{\omega}{k}$. Empirically we measure that the speed of light decreases when you enter a medium by a fixed factor, the index of refraction *n*:

 $v_{\text{medium}} = \frac{c}{n} = \frac{\omega}{k}$, so *k* must change when we enter the medium. For an air water interface we have $k_{\text{water}} = nk_{\text{air}}$. III. This is just like the harmonic waves we started with if we let $\tilde{k} = k + i\alpha$. Then we have a super compact wave: $\tilde{\psi}(x, t) = A_0 e^{i(\tilde{k}x - \omega t)}$.

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Apparently we can introduce a complex index of refraction:

$$\tilde{n} = \frac{\tilde{k}c}{\omega} = \frac{kc}{\omega} + i\frac{\alpha c}{\omega} = n + i\kappa$$
, note that $\kappa \equiv \frac{\alpha c}{\omega}$,
here κ characterizes the absorption of the wave.