

# Today

- I. Questions about Complex Waves
- II. Last Time
- III. Physical Origin of the Index of Refraction
- IV. Diffusion Equation

I. Yanpei Deng this week due to exam (will help in the lab. She's available MW from 7-8pm in Brody lab).

This week: Antu Antu will be providing homework support. Hours are: Tu 8-9pm, Th 10:30-11:30am, Th 8-9pm.

We noticed that by combining the real and imaginary exponents of our waves we could get a complex index of refraction:

$$\tilde{n} = \frac{\tilde{k}c}{\omega} = \frac{kc}{\omega} + i\frac{\alpha c}{\omega} = n + i\kappa, \text{ note that } \kappa \equiv \frac{\alpha c}{\omega}.$$

I.

$$k = \text{wave number} = \frac{2\pi}{\lambda}$$

$$\kappa = \frac{\alpha c}{\omega} = \text{imaginary part of the complex index of refraction } \tilde{n}$$

Physically this captures a property of the material that the light is going into, in particular, how much the material attenuates the wave. Specifically

$$A(x) = A(0)e^{-\alpha x}.$$

Let's focus on the real part of the complex index of refraction, the  $n$ . What is this for a material? In practice, this slows the speed of the wave:

$$v_{\text{medium}} = \frac{c}{n}$$

Suppose we had an electric field...

## II. Recall the wave equation

$$\frac{\partial^2 E(x, t)}{\partial t^2} = c^2 \frac{\partial^2 E(x, t)}{\partial x^2}$$

If you derive this wave equation from Maxwell's equations for the

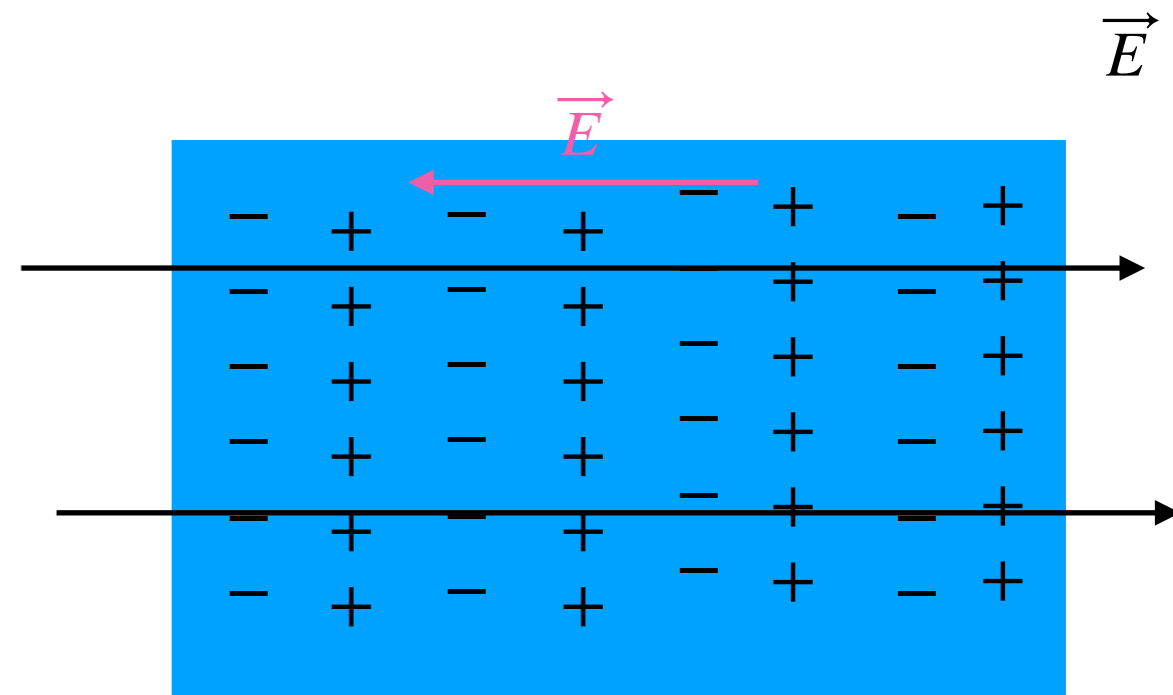
Electric and Magnetic fields you find

$$\frac{\partial^2 E(x, t)}{\partial t^2} = \frac{1}{\mu_0 \epsilon_0} \frac{\partial^2 E(x, t)}{\partial x^2}.$$

Here  $\mu_0$  is the permeability of free space and  $\epsilon_0$  is permittivity of free space. Then  $c = 1/\sqrt{\epsilon_0 \mu_0} = 3 \times 10^8 \text{ m/s}$ .

In a material we use the new permittivity:  $\epsilon_r \epsilon_0$ , here  $\epsilon_r$  is the relative permittivity and is a pure number. We also have  $\mu_r \mu_0$ , interestingly then

$$v_{\text{medium}} = \frac{1}{\sqrt{\epsilon_r \epsilon_0 \mu_r \mu_0}} = \frac{c}{\sqrt{\epsilon_r \mu_r}} \equiv \frac{c}{n}$$



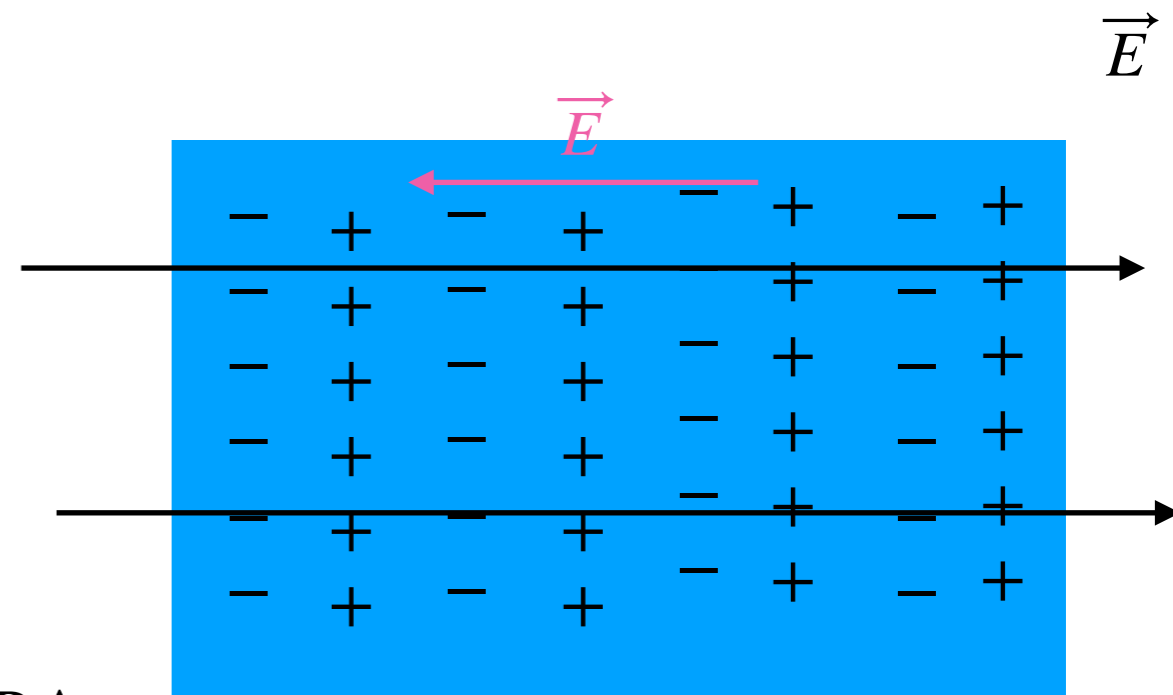
II.

For most materials we have  $\mu_r \approx 1$ . So really, the index of refraction is basically a measure of the permittivity of a medium:

$$v_{\text{medium}} = \frac{1}{\sqrt{\epsilon_r \epsilon_0 \mu_r \mu_0}} = \frac{c}{\sqrt{\epsilon_r \mu_r}} \approx \frac{c}{\sqrt{\epsilon_r}} \equiv \frac{c}{n}.$$

IV. Let's connect the random walk that you studied on the homework to the diffusion equation:

we've got a one dimensional line and because the walkers take discrete steps, they can only be located at discrete positions along the line, call the number of them at position  $i$ ,  $N_i(t)$ . Let's call the probability that a walker leaves its current location  $p = R\Delta t$



IV. Let's connect the random walk that you studied on the homework to the diffusion equation: we've got a one dimensional line and because the walkers take discrete steps, they can only be located at discrete positions along the line, call the number of them at position  $i$ ,  $N_i(t)$ . Let's call the probability that a walker leaves its current location  $p = R\Delta t$ .

Then the number walkers that go from position  $i$  to position  $i + 1$ , is  $pN_i$  in time  $\Delta t$ . In total we have

$$N_i(t + \Delta t) = N_i(t) - R\Delta t N_i(t) - R\Delta t N_i(t) + R\Delta t N_{i-1}(t) + R\Delta t N_{i+1}(t)$$