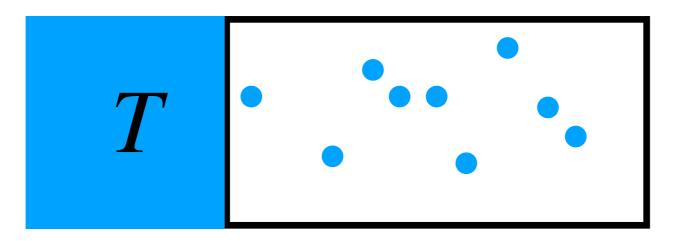
Today

- I. Last Time
- II. Spherical Coordinates and Spherical Integration
- III. Kinetic Theory of a Gas of Photons
- IV. The Power in a Gas of Photons
 - I. See below

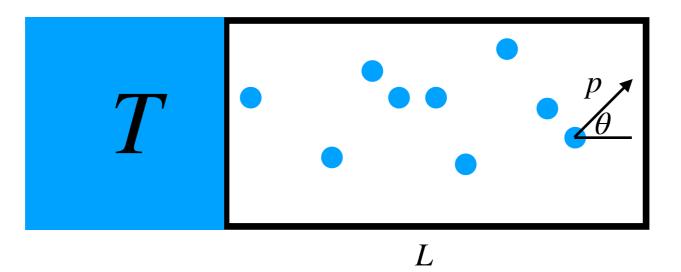
IV. Returning to the main theme of our course, let's consider a gas of photons in a box at temperature T.



"A box of light at temperature *T*". It will be convenient to introduce the energy density $u(T) = \frac{U}{V}$.

Light has energy, but also momentum. For light with energy E we have p = E/c. Since they have momentum, the light particles exert a force on the walls of the box, which leads to a pressure.

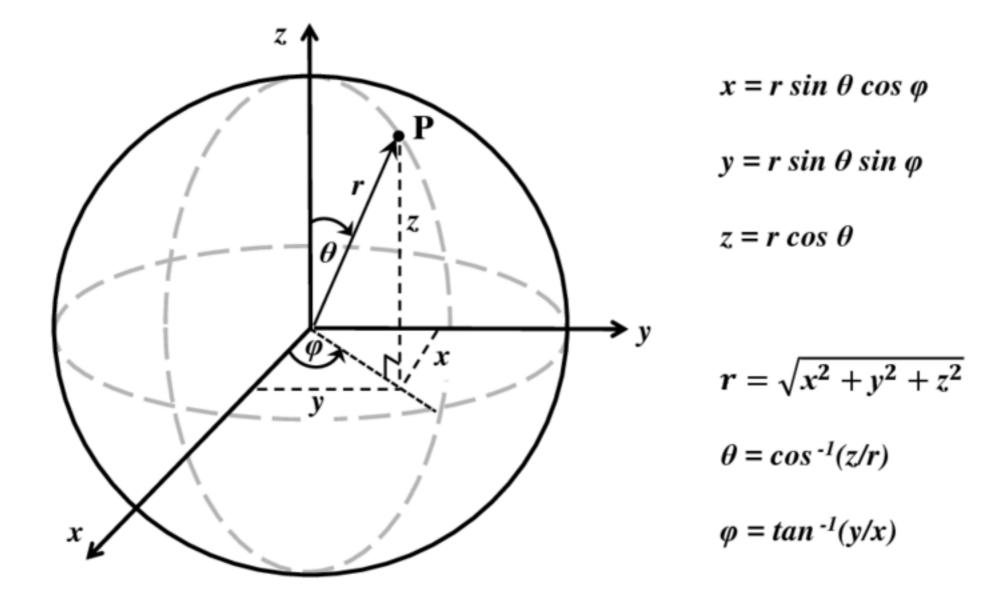
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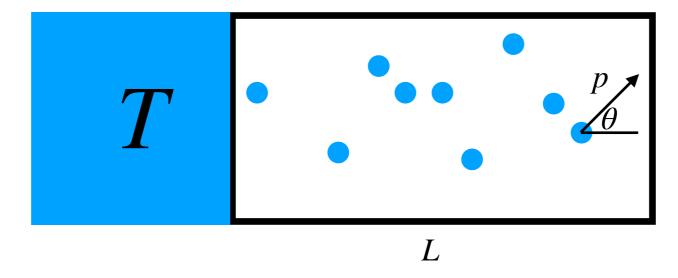
Consider a cubical box with side length *L* and hence volume $V = L^3$, then radiation at a constant temperature *T*, will also have u(T) constant. The speed of a photon at angle θ is $c \cos \theta$, and it will hit the right wall every $2L/(c \cos \theta)$ seconds.

$$F = \frac{\Delta p}{\Delta t} = \frac{2p\cos\theta}{2L/(c\cos\theta)} = \frac{pc\cos^2\theta}{L} = \frac{E\cos^2\theta}{L};$$
$$P = \frac{Force}{Area} = \frac{E\cos^2\theta/L}{L^2} = \frac{E}{V}\cos^2\theta \text{ (for 1 photon)}$$

II. Spherical Coordinates Polar angle: θ Azimuthal angles: φ, ϕ



Area Hemisphere = $\int_{0}^{2\pi} \int_{0}^{\pi/2} R^2 \sin\theta d\theta d\phi = R^2 \int_{0}^{2\pi} \left[-\cos\theta \right]_{0}^{\pi/2} d\phi = R^2 \int_{0}^{2\pi} 1 d\phi = 2\pi R^2$

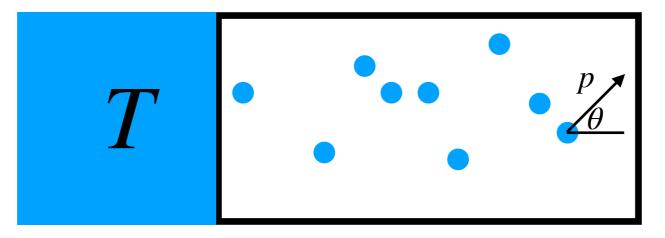


In general, we will get photons contributing in every direction

$$F = \frac{\Delta p}{\Delta t} = \frac{2p\cos\theta}{2L/(c\cos\theta)} = \frac{pc\cos^2\theta}{L} = \frac{E\cos^2\theta}{L};$$
$$P = \frac{Force}{Area} = \frac{E\cos^2\theta/L}{L^2} = \frac{E}{V}\cos^2\theta \text{ (for 1 photon)}$$

We can find the total pressure by integrating up the contributions

$$P_{tot} = \frac{1}{2\pi} \int_0^{2\pi} \int_0^{\pi/2} \frac{E}{V} \cos^2 \theta \sin \theta d\theta d\phi = \frac{E}{V} \int_0^{\pi/2} \cos^2 \theta \sin \theta d\theta d\phi$$
$$= -\frac{1}{3} \left[\cos^3 \theta \right]_0^{\pi/2} \frac{E}{V} = \frac{1}{3} \frac{E}{V$$



Switch from P_{tot} to just P and from E to U: $P = \frac{1}{3} \frac{U}{V} = \frac{1}{3} \frac{U}{V}.$

Let's compare this to what we learned for the ideal gas

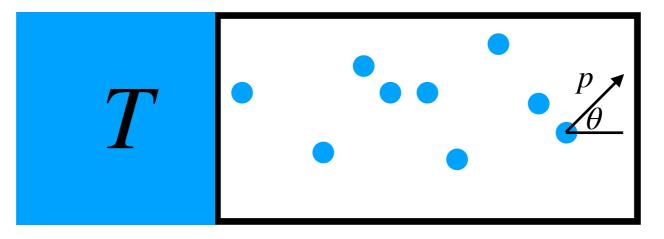
$$U = \frac{3}{2}NkT,$$
$$PV = NkT.$$

Then

$$U = \frac{3}{2}PV$$
 or $P = \frac{2}{3}\frac{U}{V} = \frac{2}{3}u$

Today I want to use without proof the Gibbs-Duhem relation:

$$Ud\left(\frac{1}{T}\right) + Vd\left(\frac{P}{T}\right) = 0.$$



Switch from P_{tot} to just P and from E to U: $P = \frac{1}{3} \frac{U}{V} = \frac{1}{3} \frac{U}{V}.$

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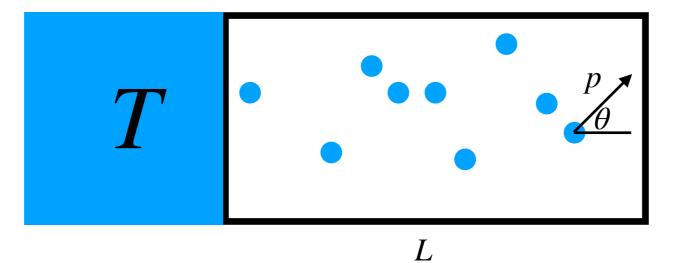
$$Ud\left(\frac{1}{T}\right) + Vd\left(\frac{P}{T}\right) = 0.$$

Divide everything by the volume

$$\frac{U}{V}d\left(\frac{1}{T}\right) + d\left(\frac{P}{T}\right) = 0 \implies 3Pd\left(\frac{1}{T}\right) + d\left(\frac{P}{T}\right) = 0,$$

Next divide both sides by P/T to get

$$\frac{3}{1/T}d\left(\frac{1}{T}\right) + \frac{1}{P/T}d\left(\frac{P}{T}\right) = 0.$$



Next divide both sides by P/T to get

$$\frac{3}{1/T}d\left(\frac{1}{T}\right) + \frac{1}{P/T}d\left(\frac{P}{T}\right) = 0.$$

Move one term to the other side to get

$$\frac{1}{P/T}d\left(\frac{P}{T}\right) = -\frac{3}{1/T}d\left(\frac{1}{T}\right)$$

And integrate both sides

$$\ln\left(\frac{P}{T}\right) = -3\ln\left(\frac{1}{T}\right) + \text{const} = \ln\left(T^3\right) + \text{const.}$$

Exponentiating both sides gives

$$\frac{P}{T} = aT^3$$
 or $P = aT^4$ or $u = \sigma T^4$.