

Today

I. Last Time

II. Wrap up Discussion of Schroedinger Equation

III. The Infinite Square Well

I. Introduced operators, which acted on the waves and returned

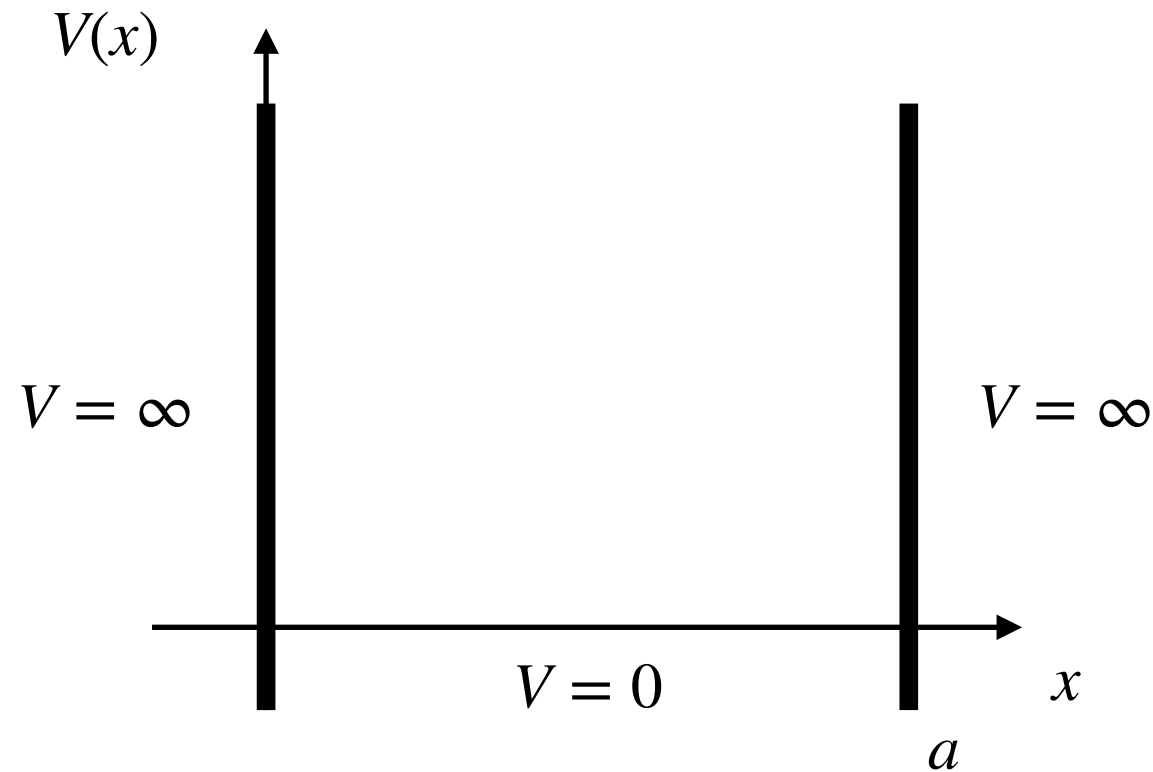
physical values: e.g. $\hat{p} = \frac{\hbar}{i} \frac{\partial}{\partial x}$ (recall $[\hbar] = \text{m kg m/s}$).

$$\hat{H}\Psi(x, t) = i\hbar \frac{\partial \Psi}{\partial t} = \frac{\hat{p}^2}{2m} \Psi + \hat{V}(x)\Psi(x, t) = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V(x)\Psi$$

This is the time-dependent Schroedinger equation. We also had

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + V(x)\psi(x) = E\psi(x) \text{ (the time-indep. Sch. Eq.)}$$

II. Today we're going to study an example called the infinite square well:



As an equation we can write this in the following form:

$$V(x) = \begin{cases} 0, & \text{if } 0 \leq x \leq a, \\ \infty, & \text{otherwise.} \end{cases}$$

For this potential we can focus our study of the Sch. Eq. to the interior of the box.

II. Where $V = 0$ we have

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi(x),$$

or

$$\frac{d^2\psi}{dx^2} = -\frac{2mE}{\hbar^2}\psi(x) = -k^2\psi(x). \quad (\text{Here } k \equiv \frac{\sqrt{2mE}}{\hbar})$$

(Recall the harmonic oscillator had Newton's equation give by

$$\frac{d^2x}{dt^2} = -\omega^2x.$$

With solutions $x(t) = A \cos(\omega t) + B \sin(\omega t)$.) Then

$$\psi(x) = A \sin(kx) + B \cos(kx). \quad (\text{check it!})$$

In addition to the Sch. Eq. , we have boundary conditions which are

$$\psi(x = 0) = 0 \quad \text{and} \quad \psi(x = a) = 0.$$

II. Then

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In addition to the Sch. Eq. , we have boundary conditions which are

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Plugging in $x = 0$ gives,

$$\psi(0) = A \sin(0) + B \cos(0) = B = 0.$$

We must have $B = 0$. Then,

$$\psi(x) = A \sin(kx).$$

We still have the boundary condition at $x = a$,

$$\psi(a) = A \sin(ka) = 0.$$

Instead of fixing $A = 0$, which is boring, we choose

$$ka = n\pi, \quad n = 0, \pm 1, \pm 2, \pm 3, \dots$$

In other words, only certain k 's work

$$k_n = \frac{n\pi}{a}, \text{ with } n = 1, 2, 3, \dots$$

II. In other words, only certain k 's work

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Physically this amounts to a requirement of only particular energies:

$$\frac{\sqrt{2mE_n}}{\hbar} = \frac{n\pi}{a}, \text{ with } n = 1, 2, 3, \dots$$

Solving for the energy we get

$$E_n = \frac{\hbar^2 k_n^2}{2m} = \frac{n^2 \pi^2 \hbar^2}{2ma^2}. \quad (n = 1, 2, 3, \dots)$$

Returning to our wave function we have

$$\psi(x) = A \sin\left(\frac{n\pi}{a}x\right), \quad n = 1, 2, 3, \dots$$

$|\psi(x)|^2 =$ probability density for particle position

