

# Today

I. Last Time

II. Degenerate Fermi Gas

III. Saiqi's Guest Lecture on Black Body Radiation

I. Last time we derived the occupancy of a Bose-Einstein gas

$$\bar{n}_{BE} = \frac{1}{e^{(\epsilon - \mu)/kT} - 1}.$$

Fermi energy:  $\epsilon_F \equiv \mu(T = 0)$ .

Let's again fix our attention on a cubical box of side length  $L$  (we're think of it as a chunk of metal). Recall that free particles have

$\lambda_n = \frac{2L}{n}$ , and  $p_n = \frac{h}{\lambda_n} = \frac{hn}{2L}$ , but this time let's take into account the

3D nature of the box  $p_x = \frac{hn_x}{2L}$ ,  $p_y = \frac{hn_y}{2L}$ ,  $p_z = \frac{hn_z}{2L}$ . The energies

$$\epsilon(\vec{n}) = \frac{|\vec{p}|^2}{2m} = \frac{h^2}{8mL^2}(n_x^2 + n_y^2 + n_z^2).$$

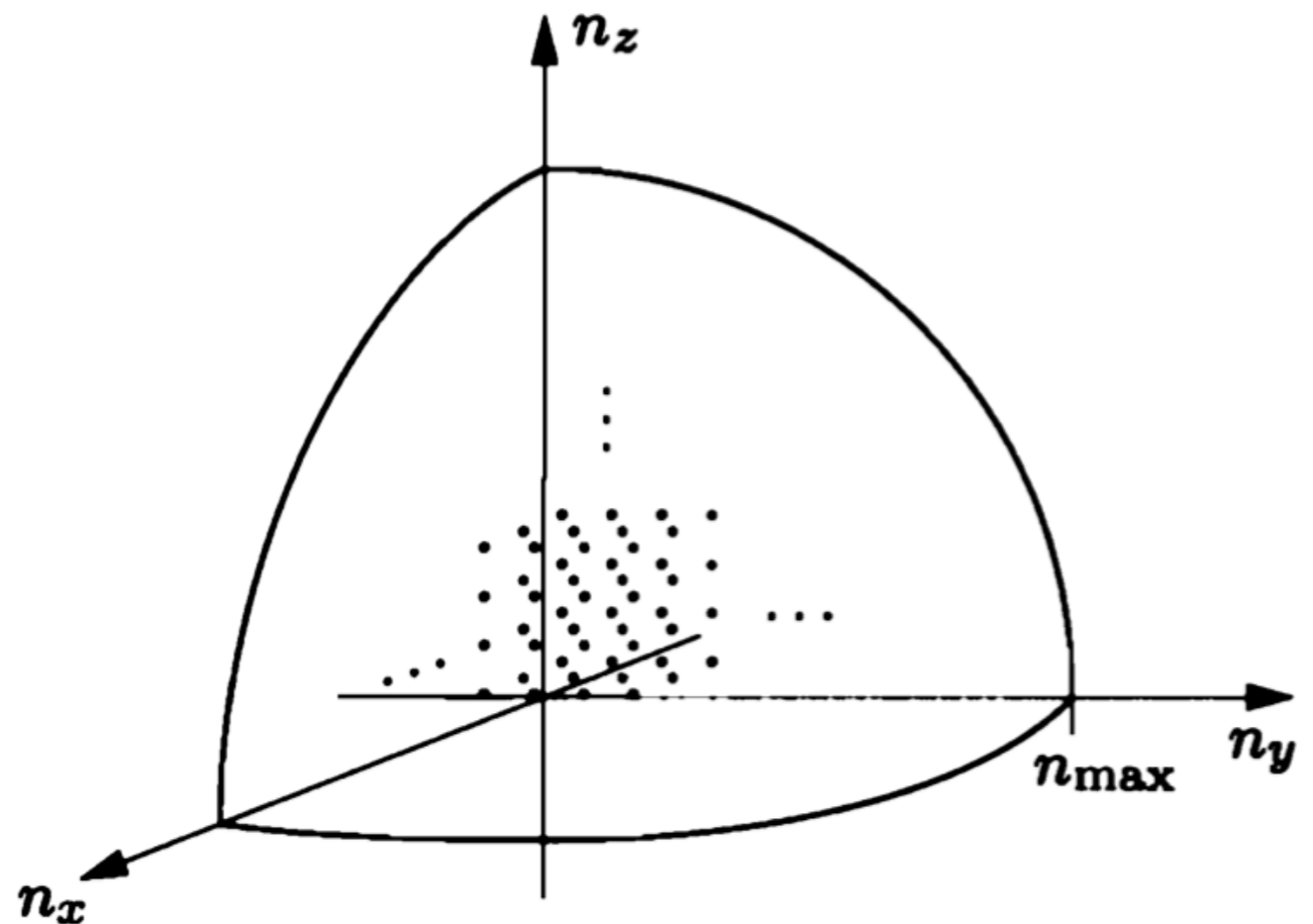
II. We'd like to compute the total number of allowed quantum states and to do that we have to count the lattice sites for positive integer  $n_x, n_y, n_z$  and a very convenient way to do that is to compute the volume of the octant of a sphere in  $n$ -space. Then we have

$$N = 2 \times (\text{vol of octant of sphere}) = 2 \cdot \frac{1}{8} \cdot \frac{4}{3} \pi n_{\text{max}}^3 ,$$

here  $n = \sqrt{n_x^2 + n_y^2 + n_z^2}$ .

On the other hand, we can solve the above for  $n_{\text{max}}$  in terms of  $N$ , and plug in to energy to get

$$\epsilon_F = \frac{h^2 n_{\text{max}}^2}{8mL^2} = \frac{h^2}{8m} \left( \frac{3N}{\pi V} \right)^{2/3}$$



II. Now that we know how far to go we can compute the  $U$ :

$$U = 2 \sum_{n_x} \sum_{n_y} \sum_{n_z} \epsilon(\vec{n}) \approx 2 \int_0^{n_{max}} \int_0^{\pi/2} \int_0^{\pi/2} n^2 \sin \theta \epsilon(n) d\phi d\theta dn = \pi \int_0^{n_{max}} \epsilon(n) n^2 dn$$

$$= \pi \int_0^{n_{max}} \frac{\hbar^2 n^2}{8mL^2} n^2 dn = \frac{\pi \hbar^2}{40mL^2} n_{max}^5 = \frac{3}{5} N \epsilon_F$$

$$n = \sqrt{n_x^2 + n_y^2 + n_z^2}$$

