

Today

- I. Announcements: One final class on Monday of completion days (Dec. 14th), Final due Friday Dec. 18th at 10pm ET
- II. Last Time
- III. Density of States Wrap Up
- IV. Balthazar's Guest Lecture on Debye Theory of Solids
- V. CAFE forms for Thermal: <https://tools.bard.edu/tools/cafeform/>

I. Bruno told us about Bose-Einstein condensation. At low temperatures all Bosons drop into the lowest energy state and they behave collectively in quantum mechanical way.

Introduced the idea of the density of states: divide energy space into bands, then the density of states tells me how many quantum states there are in a given band of energies.

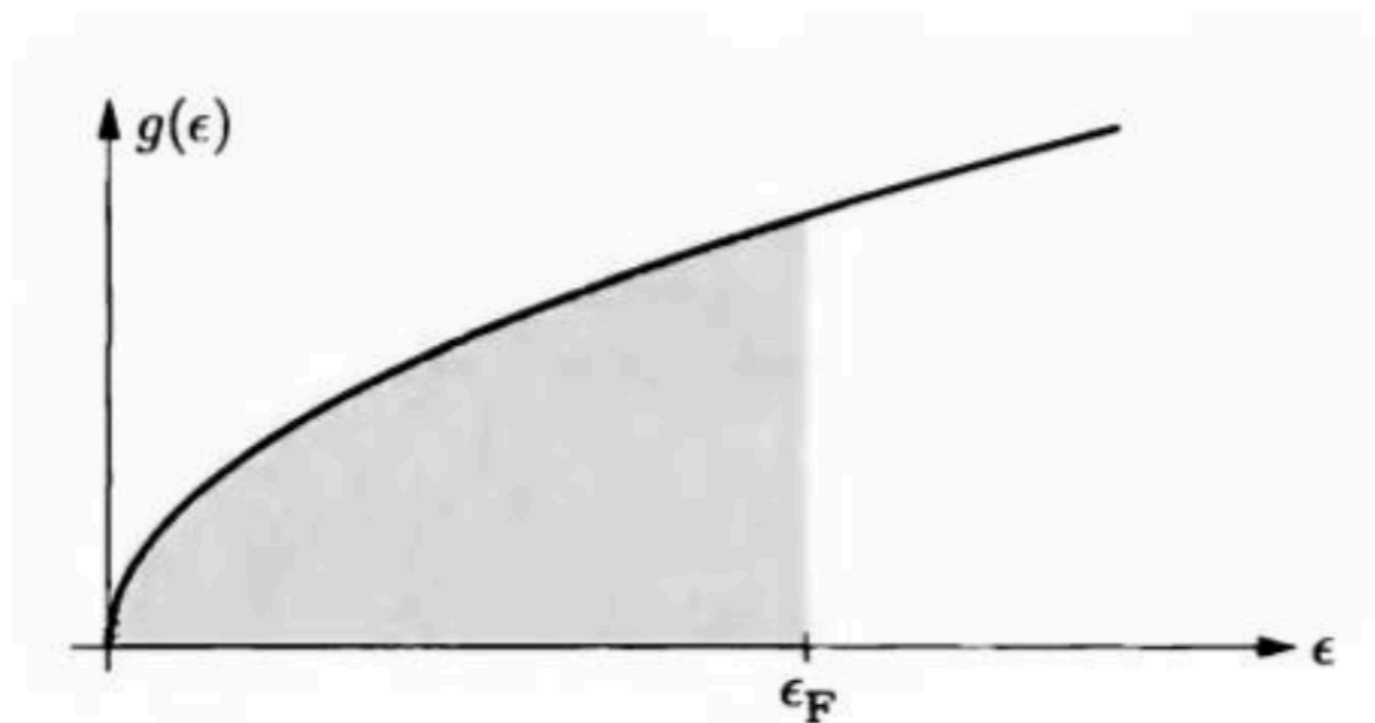
II. Density of states.

Let's recast our computation of the energy of the Fermi gas in terms of the energies of the states:

$$\epsilon = \frac{h^2}{8mL^2}n^2, \text{ gives } n = \sqrt{\frac{8mL^2}{h^2}}\sqrt{\epsilon}, \text{ and so } dn = \sqrt{\frac{8mL^2}{h^2}}\frac{1}{2\sqrt{\epsilon}}d\epsilon$$

$$U = \pi \int_0^{n_{max}} \frac{h^2 n^2}{8mL^2} n^2 dn = \int_0^{\epsilon_F} \epsilon \left[\frac{\pi}{2} \left(\frac{8mL^2}{h^2} \right)^{3/2} \sqrt{\epsilon} \right] d\epsilon$$

$\longleftarrow \hspace{10em} \hspace{10em} \longrightarrow$
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Note that we can write this in one more way:

$$g(\epsilon) = \frac{\pi(8m)^{3/2}}{2h^3} V \sqrt{\epsilon}.$$

By definition we can also compute the number of particles in our

system using this: $N = \int_0^{\epsilon_F} g(\epsilon) d\epsilon \quad (T = 0).$

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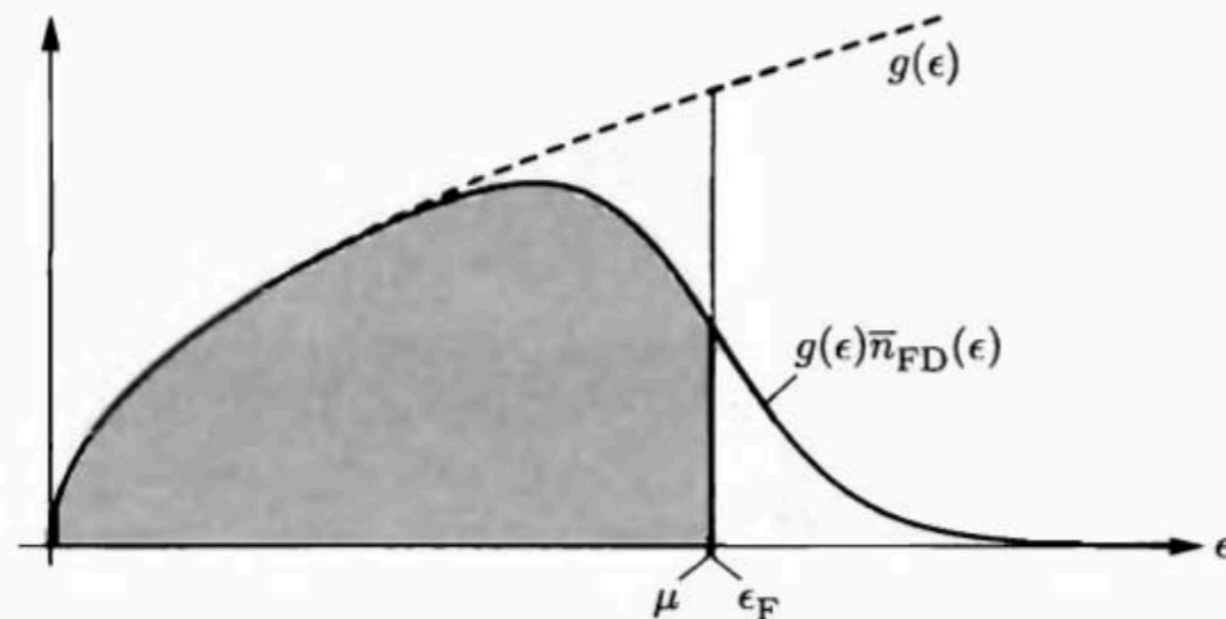
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We can do this at nonzero temperatures, as long as we take the probability of occupation into account

$$N = \int_0^{\infty} g(\epsilon) \bar{n}_{FD}(\epsilon) d\epsilon = \int_0^{\infty} g(\epsilon) \frac{1}{e^{(\epsilon-\mu)/kT} + 1} d\epsilon$$

$\mu(T) \neq \epsilon_F$ except when $T = 0$.



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