

I last time

II Fermat's principle

III Huygens' principle and wave theory

IV The opto-mechanical analogy

• Fermat's principle: Light travels the path from S to D that is extremal in time.

II last time we started to investigate an example of Fermat's principle; we considered a ray going from a source S in air to a detector D in glass.

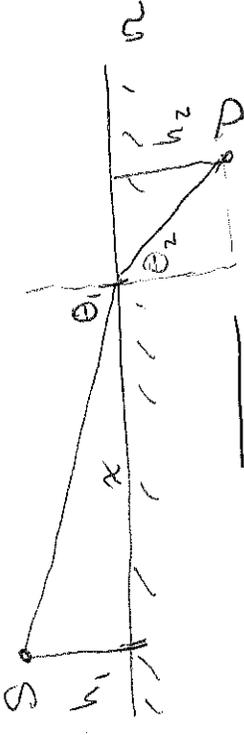
I •  $\langle f(x) \rangle = \sum_{j=1}^N f(x_j) P(x_j)$

$\sigma^2 = \langle (\Delta_j)^2 \rangle = \langle j^2 \rangle - \langle j \rangle^2$

• Probability density  $p(x)$   
 $[p(x)] = \frac{\text{Probability}}{\text{Units of } x}$  note it a density

$P(x)dx = \left\{ \begin{array}{l} \text{Prob. of event} \\ \text{between } x \text{ and } x+dx \end{array} \right\}$

of index of refraction  $n$ .



$t_1 = \frac{\sqrt{n_1^2 + x^2}}{c}$

$t_2 = \frac{\sqrt{n_2^2 + (L-x)^2}}{c/n}$

so,  $T = t_1 + t_2 = \frac{\sqrt{n_1^2 + x^2}}{c} + n \frac{\sqrt{n_2^2 + (L-x)^2}}{c}$

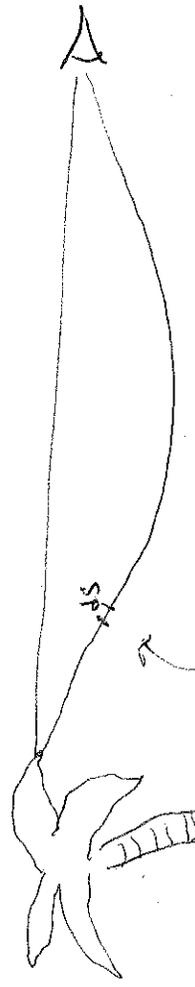
We are looking for an extremum,

$$\text{So, } \frac{dT}{dx} = 0 = \frac{1}{2} \frac{2x}{c\sqrt{h_1^2 + x^2}} - \frac{1}{2} \frac{2(L-x)}{c\sqrt{h_2^2 + (L-x)^2}}$$

$$\Rightarrow \boxed{\sin \theta_1 = n \sin \theta_2}$$

Snell's law.

Of course, the index of refraction interface need not be so simple, e.g.  $n$  could vary continuously  $n = n(h)$  different paths



This ray looks as though it's been reflected off of water.

Let's set up the general case

$$T = \int_{t_0}^{t_1} dt = \frac{c}{c} \int_{t_0}^{t_1} \frac{dt}{ds} \cdot \frac{ds}{dt} dt$$

small arc length

In fact, this explains many desert mirages — the index of refraction of air is not exactly one and varies with temperature.

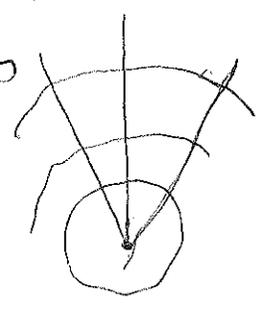
This allows light leaving a palm tree to reach your eye along two

$$= \frac{1}{c} \int_S^D \frac{c}{v} ds = \frac{1}{c} \int_S^D n ds$$

So, in general you try to find the path connecting  $S$  and  $D$  that minimizes extremizes this integral

$$\boxed{T = \frac{1}{c} \int_S^D n ds}$$

But, we've all been taught that the wave theory of light won out, and so, these rays are fictitious. Really there are wave fronts and rays are just a nice guideline overlay



of light, but let's pull in one of its beautiful principles.

If you're given a wave front, say of a plane wave, how do you construct the next one in the train?

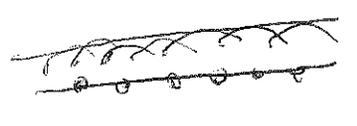
Huygen's gave an answer: break the front into a (large) collection

So, is Fermat's principle P3/4 a mathematical curiosity? A "divine miracle"? No!

Actually it follows most clearly from the wave theory.

III Huygen's principle Our goal here is not to develop in full the wave theory

of point sources, the next front is given by the superposition of spherical waves emitting from these sources



Notice the clear role of locality in this construction.

The Fermat principle follows simply from Huygens' principle. Let's take an example again.



Because the local speed changes ~~steps~~ and partially muddy field. Marchers entering the mud slow and the other marchers cover more ground, causing the wave front to bend.

III Next time we will discuss Hamilton's remarkable observation that there is a formulation

at the glass interface  $P_{y/4}$  the wave front bends. This due to wavelets in the air trap-sly catching up to those in the glass.

A useful analogy is a marching band ~~is~~ marching across a partially dry of mechanics quite similar to Fermat's formulation of optics.