

Today

General Relativity

April 13th, 2016

P1/3

I Last time

Day 29

I. The tangent space is the linear space of directional derivatives

II Tensors

- (i) Intro to definition
- (ii) Tensors are ambiguous
- (iii) Map notation

• It has a coordinate basis

$$e_\alpha = \frac{\partial}{\partial x^\alpha}$$

• A dual vector is a linear map from vectors to real numbers.

• Just like vectors, dual vectors have components. These are just the numbers that result when a dual vector eats a basis vector,

$$\omega(e_\beta) = \omega_\beta e^\beta(e_\alpha) = \omega_\beta \delta^\beta_\alpha = \omega_\alpha$$

Our definition of dual bases

• The metric allows us to identify dual vectors and vectors: "raising and lowering indices"

$a_\alpha = g_{\alpha\beta} a^\beta, \quad a^\alpha = g^{\alpha\beta} a_\beta$
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II Tensors are unnecessarily

shrouded in a fog of confusion:

"... baffling beasts bristling with indices...". The idea of

a tensor is actually quite close to what we've just described

for dual vectors. A tensor is

something that eats vectors (and dual vectors) and returns a number.

They do this in a multilinear manner, i.e. linear in each entry. In equations,

$$\begin{aligned}
 t(\underline{a}, \underline{b}, \underline{c}) &= t(a^\alpha \underline{e}_\alpha, b^\beta \underline{e}_\beta, c^\gamma \underline{e}_\gamma) \\
 &= a^\alpha b^\beta c^\gamma t(\underline{e}_\alpha, \underline{e}_\beta, \underline{e}_\gamma) \\
 &= a^\alpha b^\beta c^\gamma t_{\alpha\beta\gamma} \leftarrow \text{the components of the tensor}
 \end{aligned}$$

We call the number of vectors and dual vectors that a tensor eats its rank (= total # of indices on tensor).

Tensors are ubiquitous.

Why?! I can offer three answers:

(1) Linearity is an effective initial physical model.

Consider a spring: $F = -kx$ is an excellent approximation for small displacements in one dimension;

things get richer when you consider a 3D balloon. Name three orthogonal faces, say the x-face, y-face and z-face.

We've been spending a lot of time with a particular rank 2 tensor:

$$\begin{aligned}
 g(\underline{a}, \underline{b}) &= g(a^\alpha \underline{e}_\alpha, b^\beta \underline{e}_\beta) \\
 &= a^\alpha b^\beta g(\underline{e}_\alpha, \underline{e}_\beta) \\
 &= a^\alpha b^\beta g_{\alpha\beta} \\
 &= \underline{a} \cdot \underline{b}
 \end{aligned}$$

the metric! Claim:

If I push on the x-face a small amount, it squashes a little, and the force the balloon exerts back is, to a good first approx., linear in this displacement.

However, if I shear the x-face too this ~~shear~~ can in general change the outward force on the x-face. So, in general

$$F_x = k_{xx} X + k_{xy} Y + k_{xz} Z$$

For all three faces

$$\begin{pmatrix} F_x \\ F_y \\ F_z \end{pmatrix} = \begin{pmatrix} k_{xx} & k_{xy} & k_{xz} \\ k_{yx} & k_{yy} & k_{yz} \\ k_{zx} & k_{zy} & k_{zz} \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

or $F_i = k_{ij} x_j$

The coefficients k_{ij} are called the elasticity tensor. So, again, any time linearity is a good model tensors will show up.

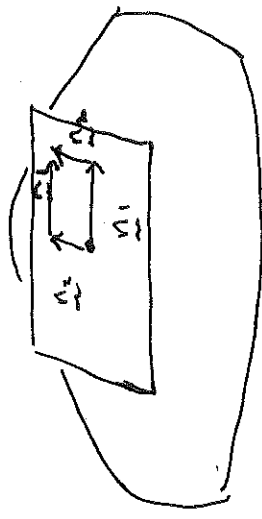
$s, e \rightarrow 0$.

(3) Tensors are great for capturing a large class of invariants. More on this shortly.

Map Notation

Call the vector space that we've been drawing vectors from V (e.g. the tangent space at P). Call its dual space V^* . Then in map notation

(2) Linearity is computationally nice. We can solve linear systems. We can do calculus with tensors:



Let s, e be small parameters, then $R(\underline{s}, \delta \underline{n}_1, \epsilon \underline{n}_2) = \epsilon \delta R(\underline{s}, \underline{n}_1, \underline{n}_2)$, and it is easy to consider the limits

$$t: V \times V \times V \rightarrow \mathbb{R}$$

$$(a, b, c) \mapsto a^{\alpha} b^{\beta} c^{\gamma} t_{\alpha\beta\gamma}$$

This notation helps us to address one of your longstanding questions: Is the metric a matrix or not?

Recall, a linear map is a map $M: V \rightarrow V$ such that M is linear TO BE CONTINUED...