

General Relativity

I Last time P1/3

Today!

II Last time

III Using Covariant Derivatives

IV Curvature Begins

April 20, 2016

Day 32

• Tensors are multilinear maps from a collection of vectors and dual vectors to a real number.

- Raise & lower tensor indices with the metric
- Defined the covariant derivative:

O Announcements

$$\nabla_{\underline{t}} \underline{v}(x^{\alpha}) = \lim_{\epsilon \rightarrow 0} \frac{[\underline{v}(x^{\alpha} + \epsilon \underline{t}^{\alpha})]_{\text{trans to } x^{\alpha}} - \underline{v}(x^{\alpha})}{\epsilon}$$

Using the key recognition that the change upon parallel transport should be proportional to v^{α} and to ϵt^{α} we found

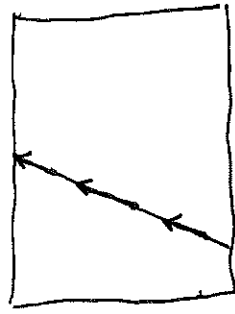
$$\nabla_{\underline{t}} v^{\alpha} = \frac{\partial v^{\alpha}}{\partial x^{\beta}} t^{\beta} + \Gamma^{\alpha}_{\beta\gamma} v^{\gamma} t^{\beta}$$

means $\underline{t} = \epsilon \underline{t}^{\beta}$

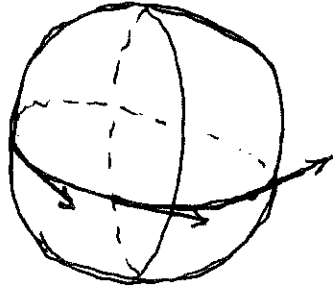
II Using Covariant Derivatives

To use the previous formula we need to identify $\Gamma^{\alpha}_{\beta\gamma}$. In addition to being extremal, geodesics are as straight as possible — meaning their tangent vectors parallel propagate into each other:

Pictorially,



or



This means that

$$(\nabla_{\tilde{u}} \tilde{u})^\alpha = 0 = u^\beta \left(\frac{\partial u^\alpha}{\partial x^\beta} + \tilde{\Gamma}^\alpha_{\beta\gamma} u^\gamma \right)$$

and comparing this with the

In special relativity we defined

$$a^\alpha = \frac{du^\alpha}{d\tau} \quad (\text{LIF only})$$

but more generally we should define acceleration as something that takes you off of geodesics, i.e.,

$$a = \nabla_{\tilde{u}} \tilde{u}$$

This is sensible because it tells you how the 4-velocity changes

geodesic equation we PZ/3 find,

$$u^\beta \left(\frac{\partial u^\alpha}{\partial x^\beta} + \Gamma^\alpha_{\beta\gamma} u^\gamma \right)$$

that $\tilde{\Gamma}^\alpha_{\beta\gamma} = \Gamma^\alpha_{\beta\gamma}$! This follows because at a point p there are geodesics in every possible direction u^γ and we can identify coefficients.

Let's do some calculations:

as you move in the direction of the 4-velocity! Also in an LIF we find,

$$\begin{aligned} (\nabla_{\tilde{u}} \tilde{u})^\alpha &= u^\beta \left(\frac{\partial u^\alpha}{\partial x^\beta} + \Gamma^\alpha_{\beta\gamma} u^\gamma \right) \quad \text{base LIF} \\ &= \frac{dx^\beta}{d\tau} \frac{\partial u^\alpha}{\partial x^\beta} = \frac{du^\alpha}{d\tau} = a^\alpha \quad \checkmark \end{aligned}$$

What acceleration is necessary to be stationary in Schwarzschild?

Well, stationary observer has,

$$(u^\alpha)^\alpha = (u^t, 0, 0, 0)$$

$$\Rightarrow -\left(1 - \frac{2M}{r}\right) (u^t)^2 = -1 \Rightarrow u^t = \left(1 - \frac{2M}{r}\right)^{-1/2}$$

and,

$$\begin{aligned} a^\alpha &= u^\beta \nabla_\beta u^\alpha = u^t \nabla_t u^\alpha = u^t \left(\frac{\partial u^\alpha}{\partial t} + \Gamma_{t\gamma}^\alpha u^\gamma \right) \\ &= u^t \left(\frac{\partial u^\alpha}{\partial t} + \Gamma_{tt}^\alpha u^t \right) \end{aligned}$$

Using the Christoffel symbols we have

$$a^\alpha = (0, \Gamma_{tt}^r u^t, 0, 0) = \left(0, \frac{M}{r^2}, 0, 0\right)$$

the other objects we've been thinking about:

$$\nabla_\alpha f = \frac{\partial f}{\partial x^\alpha}, \quad \nabla_u f = u^\alpha \frac{\partial f}{\partial x^\alpha}$$

Now use the product rule (fancier to call it the Leibniz rule):

$$\nabla_\gamma u^\alpha \omega^\beta = u^\gamma \nabla_\gamma \omega^\beta + \nabla_\gamma (u^\alpha) \omega^\beta$$

But $\nabla^\gamma \omega^\beta$ is a rank two tensor, so,

$$\nabla_\gamma t^{\alpha\beta} = \frac{\partial t^{\alpha\beta}}{\partial x^\gamma} + \Gamma_{\gamma\delta}^\alpha t^{\delta\beta} + \Gamma_{\gamma\delta}^\beta t^{\alpha\delta}$$

Interesting, it seems well behaved at $r=2M$, what's up?

Well,

$$(a \cdot a)^{1/2} = \left(1 - \frac{2M}{r}\right)^{-1/2} \frac{M}{r^2}$$

still diverges there.

Covariant Derivs of everything else:

Once you have a sensible definition for vectors, you can extend to

just like for $\nabla^\alpha \omega^\beta$.

$$\nabla_\alpha v^\beta = \frac{\partial v^\beta}{\partial x^\alpha} - \Gamma_{\alpha\gamma}^\beta v^\gamma$$

and

$$\nabla_\gamma t^{\alpha\beta} = \frac{\partial t^{\alpha\beta}}{\partial x^\gamma} + \Gamma_{\gamma\delta}^\alpha t^{\delta\beta} - \Gamma_{\gamma\delta}^\beta t^{\alpha\delta} \text{ etc}$$